

UP1289(2) : Compton Type Theory of Reflection and Refraction

The usual Compton theory can be developed as:

$$h\omega + mc^2 = h\omega_1 + \gamma mc^2 \quad (1)$$

$$h\underline{k} = h\underline{k}_1 + h\underline{k}_2 \quad (2)$$

where mc^2 is the rest energy of an electron, $h\omega$ is the energy of an incident photon, $h\omega_1$ is the energy of the scattered photon, and γmc^2 is the energy of the scattered electron.

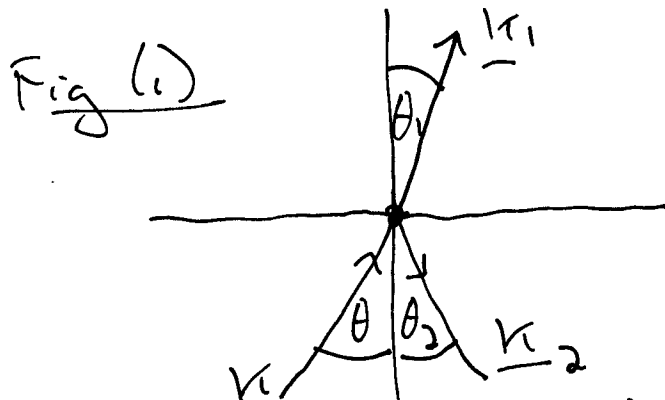


Fig (i)

With reference to Fig (i) it could be adapted directly for refraction and reflection to give the following results:

Refraction

$$\frac{1}{\omega_1} - \frac{1}{\omega} = \frac{h}{mc^2} \left(1 - \cos(\pi - \theta + \theta_1) \right) \quad (3)$$

Reflection

$$\frac{1}{\omega_2} - \frac{1}{\omega} = \frac{h}{mc^2} \left(1 - \cos(\theta + \theta_2) \right) \quad (4)$$

where m is an effective mass at the surface, and $\cos(\pi - (\theta - \theta_1)) = -\cos(\theta - \theta_1) \quad (4a)$

2) The angle between \underline{k} and \underline{k}_1 is $\frac{\pi}{2} - \theta + \theta_1$,
 and the angle between \underline{k} and \underline{k}_2 is $\theta + \theta_2$.
 By Snell's Law:

$$\theta = \theta_2 \quad - (5)$$

and

$$n \sin \theta = n_1 \sin \theta_1 \quad - (6)$$

where n is the refractive index of the medium of incidence and where n_1 is the refractive index of the medium of refraction.

Therefore:

$$\frac{1}{\omega_1} - \frac{1}{\omega} = \frac{h}{mc^2} \left(1 + \cos(\theta - \theta_1) \right) \quad - (7)$$

$$\frac{1}{\omega_2} - \frac{1}{\omega} = \frac{h}{mc^2} \left(1 - \cos(\theta + \theta_2) \right) \quad - (8)$$

where ω , ω_1 , and ω_2 are the frequencies of γ rays incident on a metal foil, then m is the mass of the electron:

$$m = 9.10953 \times 10^{-31} \text{ kg} \quad - (9)$$

At visible frequencies of about 10^{16} vibrations per second, the Compton / Moss effect is given by an effective mass of:

$m(\text{viside}) \sim 10^{-36} \text{ kg} - (10)$
 and both ω_1 and ω_2 are shifted to the red as
 is the Compton effect.

The de Broglie theory is approximately valid
 to about 100 cm^{-1} in the far infra-red, where
 $\omega = 10^{13}$ radians per second
 - (11)

So $m(\text{far infra-red}) = 10^{-38} \text{ kg} - (12)$

In the classical domain:

$$\omega = \omega_1 = \omega_2 - (13)$$

so $\cos(\theta + \theta_1) = 0 - (14)$

and $\theta + \theta_1 = \frac{\pi}{2} - (15)$

which contradicts experimental data. Similarly:

$$\sin(\theta + \theta_2) = 1 - (16)$$

and $\theta + \theta_2 = \frac{\pi}{2} - (17)$

From eqs (10) and (11):

$$\theta_1 = \theta_2 - (18)$$

Eqs (11) and (12) again contradict data.

The effective mass m is a property

The material in which reflection takes place.
 The angles θ , θ_1 and θ_2 can be related to the
 refractive index by Snell's second law (b) so
 knowing ω , ω_1 , ω_2 , θ , θ_1 and θ_2 , the effective
 mass m can be determined.

This theory regards reflection as a
 glancing collision of a moving photon from the
 effective mass m . Reflection is an acute angle
 collision in which

$$\theta = \theta_2 \quad - (19)$$

experimentally. In soft types of collision, eqs.
 (1) and (2) are obeyed.

As in HFT 158 ff Q; theory can be greatly
 developed to incorporate photon mass, and can
 also be developed for use with the Planck distribution.

By wave particle duality:

$$\frac{h}{nc^2} = \frac{1}{\omega_0} \quad - (20)$$

so:

$$\frac{1}{\omega_1} - \frac{1}{\omega} = \frac{1}{\omega_0} (1 + \cos(\theta - \theta_1))$$

$$\frac{1}{\omega_2} - \frac{1}{\omega} = \frac{1}{\omega_0} (1 - \cos(\theta + \theta_2))$$