

316(2): Vector Formulation of the JCE Identity

Consider initially the Cartan identity:

$$D_\mu T^a_{\nu\rho} + D_\rho T^a_{\mu\nu} + D_\nu T^a_{\rho\mu} := R^a_{\mu\nu\rho} + R^a_{\rho\nu\mu} + R^a_{\nu\mu\rho}$$

In UFT 254 it was shown that the vector formulation of this identity is:

$$\underline{\nabla} \cdot \underline{T}^a(\text{spin}) + \underline{\omega}^a_b \cdot \underline{T}^b(\text{spin}) = \underline{\omega}^b \cdot \underline{R}^a_b(\text{spin}) \quad \text{---(2)}$$

The JCE identity of UFT 313 is:

$$D_\mu R^a_{\lambda\rho\sigma} + D_\rho R^a_{\lambda\mu\sigma} + D_\sigma R^a_{\lambda\mu\rho} \\ := R^a_{\lambda\mu\sigma} T^d_{\rho\lambda} + R^a_{\lambda\rho\sigma} T^d_{\mu\lambda} + R^a_{\lambda\sigma\rho} T^d_{\mu\lambda} \quad \text{---(3)}$$

which can be written as:

$$D_\mu R^a_{b\rho\sigma} + D_\rho R^a_{b\mu\sigma} + D_\sigma R^a_{b\mu\rho} \\ := R^a_{b\mu\sigma} T^d_{\rho\lambda} + R^a_{b\rho\sigma} T^d_{\mu\lambda} + R^a_{b\sigma\rho} T^d_{\mu\lambda} \quad \text{---(4)}$$

By direct analogy with eq. (2) the vector formulation of eq. (4) may be written as follows, with summation over repeated c and d indices:

$$\begin{aligned}
 \underline{\nabla} \cdot \underline{R}^a_b(\text{spin}) + \underline{\omega}^a_c \cdot \underline{R}^c_b(\text{spin}) \\
 = \underline{R}^a_{bd}(\text{spin}) \cdot \underline{T}^d(\text{spin}) - (5) \\
 = \underline{R}^a_{bc}(\text{spin}) \cdot \underline{T}^c(\text{spin})
 \end{aligned}$$

This reduces to:

$$\underline{\nabla} \cdot \underline{R}(\text{spin}) + \underline{\omega}_c \cdot \underline{R}^c(\text{spin}) = \underline{R}_c(\text{spin}) \cdot \underline{T}^c(\text{spin}) - (6)$$

using: $\underline{R}(\text{spin}) = e_a e^b \underline{R}^a_b(\text{spin})$ - (7)

Multiply eq. (6) by:

$$e^c e_c = -2 - (8)$$

to obtain:

$$-2 \underline{\nabla} \cdot \underline{R}(\text{spin}) + \underline{\omega} \cdot \underline{R}(\text{spin}) = \underline{R}(\text{spin}) \cdot \underline{T}(\text{spin}) - (9)$$

So:

$$\underline{\nabla} \cdot \underline{R}(\text{spin}) = \frac{1}{2} \underline{R}(\text{spin}) \cdot (\underline{\omega} - \underline{T}(\text{spin}))$$

- (10)

Finally use:

$$3) \quad \underline{B} = W^{(0)} \underline{R}(\text{spin}) - (11)$$

to start:

$$\underline{\nabla} \cdot \underline{B} = \frac{1}{2} \underline{B} \cdot (\underline{\omega} - \underline{I}(\text{spin})) - (12)$$

The magnetic monopole is zero if and only if:

$$\underline{\omega} = \underline{I}(\text{spin}) - (13)$$

\underline{I} these equations:

$$\underline{R}^a_b(\text{spin}) = \underline{\nabla} \times \underline{\omega}^a_b - \underline{\omega}^a_c \times \underline{\omega}^c_b - (14)$$

and:

$$\underline{R}^a_b(\text{orb}) = -\underline{\nabla} \omega^a_{ob} - \frac{1}{c} \frac{d\omega^a_b}{dt} - \omega^a_{oc} \omega^c_b + \omega^a_c \omega^c_{ob} - (15)$$

with: $\underline{B}^a_b = W^{(0)} \underline{R}^a_b(\text{spin}) - (16)$

and $\underline{E}^a_b = c W^{(0)} \underline{R}^a_b(\text{orb}) - (17)$

Finally the tangent indices are removed with:

$$\underline{B} = e^a e^b \underline{B}^a_b - (18)$$

and $\underline{E} = e^a e^b \underline{E}^a_b - (19)$

It follows that:

$$4) \underline{B} = \underline{W}^{(0)} \left(\underline{\nabla} \times \underline{\omega} - \underline{\omega}_c \times \underline{\omega}^c \right) - (20)$$

$$= \underline{W}^{(0)} \underline{\nabla} \times \underline{\omega}$$

Define the magnetic flux potential by:

$$\underline{W} = \underline{W}^{(0)} \underline{\omega} - (21)$$

and

$$\underline{B} = \underline{\nabla} \times \underline{W} - (22)$$

Similarly:

$$\underline{R}(\alpha) = -\underline{\nabla} \omega_0 - \frac{1}{c} \frac{\partial \underline{\omega}}{\partial t} - \omega_{0c} \underline{\omega}^c + \underline{\omega}_c \omega_0^c - (23)$$

and

$$\underline{E} = c \underline{W}^{(0)} \left(-\underline{\nabla} \omega_0 - \frac{1}{c} \frac{\partial \underline{\omega}}{\partial t} - \omega_{0c} \underline{\omega}^c + \underline{\omega}_c \omega_0^c \right) - (24)$$
