

326(4) : Relativistic Rotational Motion

Consider the Hamiltonian and Lagrangian of special relativity:

$$H = \gamma mc^2 + U \quad (1)$$

$$L = -\frac{mc^2}{\gamma} - U \quad (2)$$

where $E = \gamma mc^2 = (p^2 c^2 + m^2 c^4)^{1/2} \quad (3)$

It follows that:

$$p^2 c^2 + m^2 c^4 = (H - U)^2 \quad (4)$$

so $(H - U)^2 - m^2 c^4 = p^2 c^2 \quad (5)$

so $H - U - mc^2 = \frac{p^2 c^2}{H - U + mc^2} \quad (6)$

Therefore:

$$H_1 = H - mc^2 = \frac{p^2 c^2}{E + mc^2} + U \quad (7)$$

In the limit

$$E \rightarrow mc^2 \quad (8)$$

$$\gamma \rightarrow 1 \quad (9)$$

i.e.

eq. (7) becomes the Schrodinger equation,

2) based on the classical Hamiltonian:

$$H_1 = \frac{p^2}{2m} + U \quad (10)$$

For rotational motion,

$$E = \frac{p^2}{2m} = \frac{1}{2} I \omega^2 = \frac{1}{2} m r^2 \omega^2 = \frac{1}{2} m v^2 \quad (11)$$

So for circular motion:

$$v = r \omega \quad (12)$$

The moment of inertia is defined as:

$$I = m r^2 \quad (13)$$

The classical kinetic energy is:

$$E = \frac{p^2}{2m} = \frac{L^2}{2I} \quad (14)$$

where
$$\underline{L} = \underline{r} \times \underline{p} \quad (15)$$

is the classical angular momentum.

The particle on a ring problem is defined by a particle moving on a circle of radius r in the xy plane. So:

$$x = r \cos \theta, \quad y = r \sin \theta \quad (16)$$

and the Laplacian is:

$$3) \quad \nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} = \frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} \quad - (17)$$

In the particle or a ring problem derivatives w.r.t. respect to r are not considered, so

$$\nabla^2 = \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} \quad - (18)$$

and the potential energy U is considered to be zero, so

$$U = 0. \quad - (19)$$

The Schrodinger equation is:

$$-\frac{1}{2m r^2} \frac{\partial^2 \psi}{\partial \theta^2} = E \psi \quad - (20)$$

The solution is the quantized rotational energy:

$$E = m_l^2 \left(\frac{\hbar^2}{2I} \right), \quad m_l = 0, \pm 1, \pm 2, \pm 3 \quad - (21)$$

(Comparing eqs. (14) and (21):

$$\boxed{L = \pm m_l \hbar} \quad - (22)$$

This is the Schrodinger quantization of rotational motion. It is similar to Sommerfeld quantization.

4) The relativistic Schrodinger equation with no potential is:

$$H_1 = H - mc^2 = \frac{p^2 c^2}{E + mc^2} \quad - (23)$$

$$= \frac{p^2 c^2}{mc^2(1+\gamma)} = \frac{p^2}{m(1+\gamma)} \quad - (23a)$$

So

$$\boxed{\frac{p^2}{m} = (1+\gamma)H_1} \quad - (24)$$

As:

$$\gamma \rightarrow 1, \quad - (25)$$

eq. (25) becomes:

$$\frac{p^2}{2m} = H_1 \quad - (26)$$

which is the free particle Schrodinger equation.

It is customary to denote:

$$H_1 := E \quad - (27)$$

Eq. (24) quantizes to:

$$\boxed{-\frac{\hbar^2}{m} \nabla^2 \psi = (1+\gamma)E\psi} \quad - (28)$$

The relativistic version of eq. (20) is:

$$-\frac{1}{2mr^2} \frac{d^2 \phi}{d\theta^2} = E_1 \phi \quad - (29)$$

where

$$E_1 = (1 + \gamma)E \quad - (30)$$

The Lorentz factor is:

$$\gamma = \left(1 - \frac{v^2}{c^2}\right)^{-1/2} \quad - (31)$$

where

$$p = mv \quad - (32)$$

so

$$\frac{p^2}{m} = \left(1 + \left(1 - \left(\frac{p}{mc}\right)^2\right)^{-1/2}\right) H_1 \quad - (33)$$

If:

$$v \ll c \quad - (34)$$

then

$$\frac{p^2}{m} = \left(1 + 1 + \frac{1}{2} \left(\frac{p}{mc}\right)^2\right) H_1 \quad - (35)$$

Using the notation of eq. (26):

$$\boxed{\frac{p^2}{m} = \left(2 + \frac{1}{2} \left(\frac{p}{mc}\right)^2\right) E} \quad - (36)$$

which is the relativistic free particle Schrödinger

b) equation in the limit (34).

Dirac quantization:

$$-\frac{\hbar^2}{m} \nabla^2 \psi = \left(2 - \frac{\hbar^2}{2mc^2} \nabla^2 \right) (E\psi) \quad (37)$$

Since E is a constant of motion:

$$\nabla^2 (E\psi) = E \nabla^2 \psi \quad (38)$$

so:

$$-\frac{\hbar^2}{m} \nabla^2 \psi = \left(2 - \frac{\hbar^2}{2mc^2} \nabla^2 \right) E \psi \quad (39)$$

i.e.

$$\left(-\frac{\hbar^2}{m} + \frac{\hbar^2 E}{2mc^2} \right) \nabla^2 \psi = 2E\psi \quad (40)$$

$$\text{or } \boxed{\left(-\frac{\hbar^2}{2m} + \frac{\hbar^2 E}{4mc^2} \right) \nabla^2 \psi = E\psi} \quad (41)$$

The Lorentz factor for the particle in a
like problem is calculated from:

$$v^2 = r^2 \dot{\theta}^2 = \left(\frac{L_0}{mr} \right)^2 \quad (42)$$

where L_0 is the non-relativistic angular
momentum defined by:

$$L_0 = m r^2 \dot{\theta} \quad (43)$$

so

$$\gamma = \left(1 - \left(\frac{L_0}{m r c} \right)^2 \right)^{-1/2} \quad (44)$$

$$\approx 2 + \frac{1}{2} \left(\frac{L_0}{m r c} \right)^2 \quad (44a)$$

where

$$v \ll c \quad (45)$$

so

$$E_1 = \left(2 + \frac{1}{2} \left(\frac{L_0}{m r c} \right)^2 \right) \quad (46)$$

in eq. (29). The expectation value of E_1 is shifted to

$$\langle E_1 \rangle = \int \psi^* E_1 \psi \, d\tau \quad (47)$$