

Note 333(1) : Classif-ctia Scheme for the New Hyperfine structure.

The classif-ctia scheme is based on the following four types of $SU(2)$ Hamiltonian in special relativity.

Class 1 $H_0(1) = \frac{1}{m} \frac{\sigma \cdot p_0}{1+\gamma} \frac{\sigma \cdot p_0}{1+\gamma} + U - (1)$

Class 2 $H_0(2) = \frac{\gamma}{m} \frac{\sigma \cdot p_0}{1+\gamma} \frac{\sigma \cdot p_0}{1+\gamma} + U - (2)$

Class 3 $H_0(3) = \frac{\gamma^2}{m} \frac{\sigma \cdot p_0}{1+\gamma} \frac{\sigma \cdot p_0}{1+\gamma} + U - (3)$

Class 4 $H_0(4) = \frac{1}{m} \frac{\gamma^2}{1+\gamma} \frac{\sigma \cdot p_0}{1+\gamma} \frac{\sigma \cdot p_0}{1+\gamma} + U - (4)$

also $\gamma = \left(1 - \frac{p_0^2}{m^2 c^2} \right)^{-1/2} - (5)$

The relativistic momentum is:

$$\underline{p} = \gamma \underline{p}_0 - (6)$$

In the presence of a magnetic field:

$$\underline{p}_0 \rightarrow \underline{p}_0 - e\underline{A} - (7)$$

In ECE and ECE2 these equations can be derived from Cartesian geometry, and \underline{A} related

2) to quantify as in UFT 317, and UFT 318.

The four classes of equations give four different types of hyperfine structure. In the presence of a magnetic field, a whole new effect appears, notably in ESR and NMR.

Eqs. (1) to (4) derive from the Hamiltonian of special relativity

$$H = E + U \quad (8)$$

where

$$E^2 = c^2 p^2 + m^2 c^4 \quad (9)$$

in ECE and ECE2 & further energy eqn (5) derived from the tetrad postulate as in previous work. From eqs. (8) and (9):

$$(H - U)^2 - m^2 c^4 = c^2 p^2 \quad (10)$$

$$\text{so } H - U - mc^2 = \frac{c^2 p^2}{H - U + mc^2} \quad (11)$$

$$\begin{aligned} \text{i.e. } H_0 = H - mc^2 &= \frac{c^2 p^2}{E + mc^2} + U \\ &= \frac{p^2}{m(1 + \gamma)} + U \quad (12) \\ &= \frac{\gamma^2 p_0^2}{m(1 + \gamma)} + U \end{aligned}$$

QED. Eqs. (1) to (4) are ways of expressing eq. (12) in the $SU(2)$ basis.

3) The old approximation by Dirac or one of his contemporaries

is:

$$H \sim mc^2 - (13)$$

but as shown in UFT 322 eq. (13) leads to:

$$H_0 = ? \cdot 0 - (14)$$

which is very restrictive if not unphysical. This fact was first discovered in UFT 322.

There is no way of knowing a priori which scheme (1) to (4) should be used. They must be carefully investigated and systematically evaluated, their new predictions compared with experimental data. This is a rigorous new test of relativistic quantum mechanics, especially in the areas of ESR and NMR.

Clearly the Dirac approximation (13) is very rough, because the correct Hamiltonian is:

$$H = E + U = \gamma mc^2 + U - (15)$$

so Dirac assumed:

$$\gamma \rightarrow 1 - (16)$$

$$U \ll mc^2 - (17)$$

and

In the H atom:

$$U = -\frac{e^2}{4\pi\epsilon_0 r} - (18)$$

4) Therefore the magnitude of U is

$$|U| = \frac{e^2}{4\pi\epsilon_0 r} \quad - (19)$$

So eq. (17) means:

$$\frac{e^2}{4\pi\epsilon_0 r} \ll mc^2 \quad - (20)$$

i.e.

$$r \gg \frac{e^2}{4\pi\epsilon_0 mc^2} \quad - (21)$$

Note that

$$\frac{e^2}{4\pi\epsilon_0 mc^2} = \frac{\hbar}{mc} \alpha \quad - (21)$$

where the fine structure constant is:

$$\alpha = \frac{e^2}{4\pi\hbar c \epsilon_0} = 0.00729351 \quad - (22)$$

and the Compton wavelength is

$$\lambda_c = \frac{\hbar}{mc} = 3.8616 \times 10^{-13} \text{ m} \quad - (23)$$

So Dirac assumed:

$$r \gg 2.816 \times 10^{-15} \quad - (24)$$

and

$$v_0 \ll c \quad - (25)$$

The Bohr radius is

$$a_0 = 5.29177 \times 10^{-11} \text{ m} \quad - (26)$$

The restrictive nature of eq. (15) only becomes clear when it is realized that

$$\gamma mc^2 + U \sim mc^2 \quad (27)$$

so

$$\gamma \sim 1 - \frac{U}{mc^2} \quad (28)$$

The correct γ is given by eq. (5), and the same approximation is used by Dirac,

$$p_0^2 \ll m^2 c^2 \quad (29)$$

$$\gamma \sim 1 + \frac{1}{2} \frac{p_0^2}{m^2 c^2} \quad (30)$$

From eqs. (28) and (30):

$$\frac{p_0^2}{2m} = -U \quad (31)$$

so

$$H_0 = \frac{p_0^2}{2m} + U = ? \quad 0 \quad (32)$$

This makes it clear that the Dirac approximation is at best very restrictive, at worst unphysical.

So it is replaced by eqs. (1) to (4)
