

533(b) : Condition to Fire Structure For any Velocity

Consider the Lagrangian:

$$H = \frac{\sigma \cdot p_0}{1+\gamma} \frac{\gamma^2}{1+\gamma} \frac{\sigma \cdot p_0}{1+\gamma} + U \quad - (1)$$

By computer algebra:

$$\nabla \left(\frac{\gamma^2}{1+\gamma} \right) = - \left[\frac{2 + \left(\frac{1-p_0^2}{m^2 c^2} \right)^{-1/2}}{\left(\left(\frac{1-p_0^2}{m^2 c^2} \right)^{1/2} + \left(\frac{1-p_0^2}{m^2 c^2} \right)^{1/2} \right)} \right] \frac{e^2}{4\pi \epsilon_0 m c^2 r^3} \quad - (2)$$

where no limits used:

$$p_0^2 = 2(H_0 - U) \quad - (3)$$

and

$$U = \frac{-e^2}{4\pi \epsilon_0 r} \quad - (4)$$

Define:

$$A := \frac{2 + \left(\frac{1-p_0^2}{m^2 c^2} \right)^{-1/2}}{\left(\left(\frac{1-p_0^2}{m^2 c^2} \right)^{1/2} + \left(\frac{1-p_0^2}{m^2 c^2} \right)^{1/2} \right)} \quad - (5)$$

Now quantize eq. (1):

$$H\psi = -i\hbar \underline{\sigma} \cdot \underline{\nabla} \left(\frac{\gamma^2}{1+\gamma} \underline{\sigma} \cdot \underline{p}_0 \psi \right) + U\psi \quad (6)$$

$$= -i\hbar \underline{\sigma} \cdot \underline{\nabla} \left(\frac{\gamma^2}{1+\gamma} \right) \underline{\sigma} \cdot \underline{p}_0 \psi + \dots + U\psi$$

From eqns. (2) and (6):

$$H\psi = \frac{ie^2\hbar}{4\pi\epsilon_0 m^2 c^2 r^3} A \underline{\sigma} \cdot \underline{r} \underline{\sigma} \cdot \underline{p}_0 \psi + U\psi \quad (7)$$

By Pauli algebra:

$$\underline{\sigma} \cdot \underline{\nabla} \underline{\sigma} \cdot \underline{p}_0 = \underline{\nabla} \cdot \underline{p}_0 + i \underline{r} \times \underline{p}_0 \quad (8)$$

where the orbital angular momentum is:

$$\underline{L} = \underline{r} \times \underline{p}_0 \quad (9)$$

So the spin orbit Hamiltonian is:

$$H_{so} = -\frac{e^2\hbar}{4\pi\epsilon_0 m^2 c^2 r^3} A \underline{\sigma} \cdot \underline{L} \psi \quad (10)$$

giving:

$$E = \langle H_{so} \rangle = \frac{-e^2\hbar A}{4\pi\epsilon_0 m^2 c^2} \left\langle \frac{\underline{\sigma} \cdot \underline{L}}{r^3} \right\rangle \quad (11)$$

It is possible to vary p_0 for zero to

infinity and plot the shift in the spin orbit fine structure, or use

$$\frac{P_0^2}{m^2 c^2} = \left\langle \frac{P_0^2}{m^2 c^2} \right\rangle = \frac{\lambda c}{a_0} \frac{d}{n^3} - (12)$$
$$= \frac{2.662567 \times 10^{-5}}{n^3}$$

and plot the resulting spectrum.

So this is the first time that the theory of spin orbit coupling has been worked out for relativistic velocities.
