

425(1): Self Consistent Results from the Lagrangian and Hamiltonian Methods

The Lagrangian of a theory can be written as:

$$L = -mc^2 \left(m(r_1) - \frac{1}{c^2} (\dot{r}_1^2 + r_1^2 \dot{\phi}^2) \right)^{1/2} + \frac{nm\hbar}{r_1} \quad (1)$$

also $v_1^2 = \dot{r}_1^2 + r_1^2 \dot{\phi}^2 \quad (2)$
is the coordinate system (r_1, ϕ) . In a central system:

$$v_1 = \dot{r}_1 \quad (3)$$

and the Euler Lagrange equation is:

$$\frac{d}{dt} \frac{\partial L}{\partial v_1} = \frac{\partial L}{\partial r_1} \quad (4)$$

It follows that:

$$\frac{d}{dt} (\gamma m v_1) = -\frac{mc^2}{2} \gamma \frac{dm(r_1)}{dr_1} - \frac{nm\hbar}{r_1^2} \quad (5)$$

The magnitude of the vacuum force is:

$$F(\text{vac}) = -\frac{mc^2}{2} \gamma \frac{dm(r_1)}{dr_1} \quad (6)$$

also
$$\gamma = \left(m(r_1) - \frac{v_1^2}{c^2} \right)^{-1/2} \quad (7)$$

The Euler Lagrange equation:

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{\phi}} = \frac{\partial L}{\partial \phi} \quad (8)$$

gives the second Euler-Lagrange Equation:

$$2) \quad \frac{dL}{dt} = 0 \quad - (9)$$

where

$$L = \gamma m(r) r_1^2 \dot{\phi} \quad - (10)$$

is the angular momentum.

The Hamiltonian is:

$$H = m(r_1) \gamma mc^2 - \frac{mM\Gamma}{r_1} \quad - (11)$$

and is related to the Lagrangian:

$$L = -\frac{mc^2}{\gamma} + \frac{mM\Gamma}{r_1} \quad - (12)$$

by

$$H = \frac{p_1^2 c^2}{\gamma mc^2} - L \quad - (13)$$

and the first principles of Lagrangian and Hamiltonian dynamics. For self consistency, eqs. (11) to (13) shows that:

$$\frac{dm(r_1)}{dt} = 0 \quad - (14)$$

If Hamiltonian is a constant of motion, so:

$$\frac{dH}{dt} = 0 \quad - (15)$$

It follows that:

$$\begin{aligned} mc^2 \frac{d}{dt} (m(r_1) \gamma) &= \frac{d}{dt} \left(\frac{mM\Gamma}{r_1} \right) \quad - (16) \\ &= -\frac{mM\Gamma}{r_1^2} \frac{dr_1}{dt} = -\frac{mM\Gamma}{r_1^2} v_1 \end{aligned}$$

Using eq. (14): -(17)

$$m c^2 n(r_1) \frac{dY}{dt} = -\frac{m M G}{r_1^2} v_1 = m c^2 n(r_1) \frac{dY}{dv_1} \frac{dv_1}{dt}$$

Now we:
$$\frac{dY}{dv_1} = Y^3 \frac{v_1}{c^2} \quad -(18)$$

t. find that:

$$m n(r_1) Y^3 \frac{dv_1}{dt} \cdot v_1 = -\frac{m M G}{r_1^2} v_1 \quad -(19)$$

so

$$m(r_1) Y^3 \frac{dv_1}{dt} = -\frac{M G}{r_1^2} \quad -(20)$$

Eq. (20) is based on:

$$\frac{dm(r_1)}{dt} = 0 \quad -(21)$$

and

$$\frac{dn(r_1)}{dv_1} = 0 \quad -(22)$$

Eq. (21) and (22) also follow from the fact that the infinitesimal line element:

$$ds^2 = c^2 dt^2 = m(r) c^2 dt^2 - \frac{dr^2}{m(r)} - r^2 d\phi^2 \quad -(23)$$

a steady state line element. Combining eqs. (19) and (20):

$$4) \quad \frac{d}{dt} (\gamma m v_i) = m m(r_i) \gamma^3 \frac{dv_i}{dt} - \frac{m c^2}{2} \gamma \frac{dm(r_i)}{dr_i} \quad -(24)$$

In special relativity Eq. equation reduces to:

$$\frac{d}{dt} (\gamma v) = \gamma^3 \frac{dv}{dt} \quad -(25)$$

which is a well known equation of special relativity, Q.E.D. Eq. (25) is proven as follows:

$$\begin{aligned} \frac{d}{dt} (\gamma v) &= \frac{d}{dv} (\gamma v) \frac{dv}{dt} = \left(\gamma + v \frac{d\gamma}{dv} \right) \frac{dv}{dt} \\ &= \left(\gamma + \frac{v}{c^2} \gamma^3 \right) \frac{dv}{dt} \\ &= \gamma^3 \frac{dv}{dt} \left(\frac{1}{\gamma^2} + \frac{v^2}{c^2} \right) \\ &= \gamma^3 \frac{dv}{dt} \left(1 - \frac{v^2}{c^2} + \frac{v^2}{c^2} \right) \\ &= \gamma^3 \frac{dv}{dt} \quad -(26) \end{aligned}$$

Q.E.D. In order to transform eq. (24) to the plane polar system (r_i, ϕ) use eq. (2)