

430(3): The Casimir Force in n Theory

The Casimir force is given by eq. (19) of the previous note:

$$\langle F \rangle = \frac{\hbar^2}{2m} \int \psi^* \nabla^2 \left(\frac{dm(r)}{dr} \cdot \frac{1}{\left(2m(r) - r \frac{dm(r)}{dr} \right)} \right) \psi d\tau \quad (1)$$

where m is an effective mass and where ψ is the wavefunction of the Casimir effect:

$$\psi_n(x, y, z, t) = \exp(-i(\omega_n t - k_x x - k_y y)) \times \sin(k_n z) \quad (2)$$

given in: www.cs.mcgill.ca/~rvest/ - (3)

Here k_x and k_y are wave numbers parallel to the capacitor plate of the Casimir effect. The wave numbers are defined by:

$$k_n = \frac{n\pi}{r} \quad (4)$$

Let

$$f(r) = \nabla^2 \left(\frac{dm(r)}{dr} \cdot \frac{1}{\left(2m(r) - r \frac{dm(r)}{dr} \right)} \right) \quad (5)$$

then:

$$\langle F \rangle = \frac{\hbar^2}{2m} \int \psi^* f(r) \psi d\tau \quad (6)$$

$$= \frac{\hbar^2}{2m} \int f(r) \psi^* \psi d\tau$$

From eq. (2):

$$\psi_n^* \psi_n = \sin^2(k_n z) \quad (7)$$

so the Casimir force for mode n is:

$$\langle F_n \rangle = \frac{\hbar^2}{2m} \int f(r) \sin^2(k_n z) d\tau \quad - (8)$$

$$= \frac{\hbar^2}{2m} \int f(r) \sin^2\left(\frac{n\pi z}{r}\right) d\tau$$

the total Casimir force for all n modes is:

$$\langle F \rangle = \sum_n \langle F_n \rangle \quad - (9)$$

$$= \frac{\hbar^2}{2m} \int f(r) \sum_n \sin^2\left(\frac{n\pi z}{r}\right) d\tau$$

and the Casimir force per unit area A is:

$$\frac{\langle F \rangle}{A} = \frac{\hbar^2}{2mA} \int f(r) \sum_n \sin^2\left(\frac{n\pi z}{r}\right) d\tau \quad - (10)$$

Here A is the area of the plates between which the Casimir force is measured.

If there is only one mode of radiation, i.e. monochromatic radiation, then:

$$n = 1 \quad - (11)$$

$$\text{and } \frac{\langle F \rangle}{A} = \frac{\hbar^2}{2mA} \int f(r) \sin^2\left(\frac{\pi z}{r}\right) d\tau \quad - (12)$$

In spherical polar coordinates:

$$\cos \phi = \frac{z}{r} \quad - (13)$$

so the Casimir force per area is:

$$3) \quad \frac{\langle F \rangle}{A} = \frac{\hbar^2}{2mA} \int f(r) \sin^2(n\pi \cos\phi) dr - (14)$$

For n modes of radiation:

$$\frac{\langle F \rangle}{A} = \frac{\hbar^2}{2mA} \int f(r) \sum_n \sin^2(n\pi \cos\phi) dr - (15)$$

This is a hybrid theory that still retains elements of quantum electrodynamics for wavefunctions (2).
The Casimir effect can be developed on the classical level using eq. (39) of UFT 428:

$$H = \frac{1}{m(r)^{1/2}} \frac{p^2}{2m} \left(1 + \frac{U_0}{2mc^2} \right) + m(r)^{1/2} (U_0 + mc^2) - (16)$$

$$= E + U$$

where the total relativistic energy is:

$$E = \frac{1}{m(r)^{1/2}} \frac{p^2}{2m} \left(1 + \frac{U_0}{2mc^2} \right) + m(r)^{1/2} mc^2 - (17)$$

and $U = m(r)^{1/2} U_0 - (18)$

$$E_0 = m(r)^{1/2} mc^2 - (19)$$

Here

is the rest energy, and

$$U_0 = \frac{-e^2}{4\pi\epsilon_0 r} - (20)$$

4) is the Coulombic potential between two charges.
 The Casimir force can be considered classically
 as eq. (14) of UFT 430(2):

$$F = - \frac{dm(r)}{dr} \left(\frac{m(r)^{1/2}}{2m(r) - r \frac{dm(r)}{dr}} \right) E \quad (21)$$

where E is defined in general by $\leftarrow r$ (17). The relevant E is obtained by removing the rest energy from consideration, so:

$$E_0 = E - m(r)^{1/2} mc^2 \quad (22)$$

and in the limit:

$$U_0 \ll 2mc^2 \quad (23)$$

we obtain:

$$E_0 = \frac{1}{m(r)^{1/2}} \frac{p^2}{2m} \quad (24)$$

The Casimir force is therefore:

$$F_0 = - \frac{dm(r)}{dr} \left(\frac{m(r)^{1/2}}{2m(r) - r \frac{dm(r)}{dr}} \right) E_0$$

$$= - \frac{dm(r)}{dr} \left(\frac{m(r)^{1/2}}{2m(r) - r \frac{dm(r)}{dr}} \right)^{-1}$$

$$F_0 = - \frac{dm(r)}{dr} \frac{p^2 / (2m)}{2m(r) - r \frac{dm(r)}{dr}} \quad (26)$$

5) Eq. (26) is an entirely new type of force in physics, it is a force between the particle of mass m and a capacitor plate, and is neither gravitational nor electromagnetic. It is due purely to the nature of n space. It is an attractive force as measured experimentally using two very accurately engineered plates.

The units of the force in eq. (26) are:

$$F_0 = \text{kg m s}^{-2} / \text{m} = \text{kg m s}^{-2} \quad \checkmark \checkmark \quad (27)$$

The force can be expressed as:

$$F_0 = -f(r) T \quad (28)$$

also

$$f(r) = \frac{dm(r)/dr}{2m(r) - r \frac{dm(r)}{dr}} \quad (29)$$

and

$$T = \frac{1}{2} m v^2 \quad (30)$$

Eq. (28) is the first classical description of the Casimir effect, and is a clear demonstration of the ability of n space to produce an experimentally observable force the well known Casimir force. In an n space with property:

$$2n(r) = r \frac{dn(r)}{dr} \quad (31)$$

The Casimir force is amplified to infinity.