

### SOME NOTES ON THE INVARIANCE OF B CYCLICS

#### First Example

Initially:

$$B^{(3)*} = -\frac{i}{B^{(0)}} B^{(1)} \times B^{(2)} \quad (1)$$

Consider a boost in the Z direction, then:

$$B^{(3)} \rightarrow B^{(3)} \quad (2)$$

$$B^{(1)} \times B^{(2)} \rightarrow \gamma^2 \left[ \begin{array}{l} B^{(1)} \times B^{(2)} - \frac{1}{c^2} (\mathbf{v} \times \mathbf{E}^{(1)}) \times B^{(2)} \\ -\frac{1}{c^2} B^{(1)} \times (\mathbf{v} \times \mathbf{E}^{(2)}) + \frac{1}{c^4} (\mathbf{v} \times \mathbf{E}^{(1)}) \times (\mathbf{v} \times \mathbf{E}^{(2)}) \end{array} \right] \quad (3)$$

The left hand side is the  $\kappa$  frame and the right hand side is the  $\kappa'$  frame, moving with respect to  $\kappa$  at  $v$ . However, we know that the same boost produces eqn. (2). So the two results are compatible if and only if  $v = c$ . We then obtain eqn. (1) again, (re. Enigmatic Photon, Vol. 4, p. 91). However, such a boost means that an observer on the photon will record its velocity as zero. This is compatible with special relativity.

#### **Conclusion**

A Lorentz boost applied to the B Cyclic theorem is not physically meaningful in any direction. The B Cyclic theorem is invariant under Lorentz transformation.

#### Second Example

Consider the plane wave in vacuo in frame  $\kappa$ :

$$E = \frac{E^{(0)}}{\sqrt{2}} (i - ij) e^{i\phi} \quad (1)$$

Then in frame  $\kappa'$ :

$$E' = \frac{\gamma}{\sqrt{2}} \left[ (E^{(0)} - vB^{(0)}) \mathbf{i} + i(E^{(0)} + vB^{(0)}) \mathbf{j} \right] e^{i\phi} \quad (2)$$

where  $v$  is the velocity of frame  $\kappa'$  with respect to  $\kappa$ . However, we know that in vacuo:

$$E^{(0)} = cB^{(0)} \quad (3)$$

So eqn. (2) becomes:

$$\begin{aligned} E' &= \frac{E^{(0)}}{\sqrt{2}} \begin{pmatrix} 0 \\ 0 \end{pmatrix} (i - ij) e^{i\phi} \\ &= E \end{aligned} \quad (4)$$

The Maxwell's equations in vacuo are invariant under Lorentz transform, which cannot be applied to them.