

174(7): Complete Solution of the Dirac Equation for the H Atom

Consider the Dirac equation in the form:

$$(E + c \underline{\sigma} \cdot \underline{\hat{p}}) \phi^L = mc^2 \phi^R \quad (1)$$

$$(E - c \underline{\sigma} \cdot \underline{\hat{p}}) \phi^R = mc^2 \phi^L \quad (2)$$

$$(E - c \underline{\sigma} \cdot \underline{\hat{p}}) \phi^L = mc^2 \phi^R \quad (3)$$

given by the Lorentz transforms of ϕ^L and ϕ^R . Use the transformations:

$$\phi^L = \phi_S^R + \phi_S^L \quad (3)$$

$$\phi^R = \phi_S^R - \phi_S^L \quad (4)$$

to find:

$$(E - mc^2) \phi_S^R + c \underline{\sigma} \cdot \underline{\hat{p}} \phi_S^L = 0 \quad (5)$$

$$(E + mc^2) \phi_S^L + c \underline{\sigma} \cdot \underline{\hat{p}} \phi_S^R = 0 \quad (6)$$

For the H atom the solutions are:

$$\phi_S^R = F |j m l - \rangle \quad (7)$$

$$\phi_S^L = -i f |j m l + \rangle \quad (8)$$

In these equations:

$$\underline{\sigma} \cdot \underline{\hat{p}} = \underline{\sigma} \cdot \underline{\hat{r}} \left(-i \hbar \frac{\partial}{\partial r} + i \underline{\sigma} \cdot \underline{\hat{L}} \right) \quad (9)$$

with

$$\underline{\sigma} \cdot \underline{\hat{r}} |j m l_{\pm} \rangle = - |j m l_{\pm} \rangle \quad (10)$$

and

$$\underline{\sigma} \cdot \underline{\hat{L}} |j m l \rangle = \hbar (j(j+1) - l(l+1) - s(s+1)) |j m l \rangle \quad (11)$$

where

$$l_{\pm} = j \pm \frac{1}{2} \quad (12)$$

The structure of the H atom is in fact given completely by eq. (11); so the atomic structure is given

2) immediately by the same equation, fine detail included. This fact is emphasized using equation (9), which gives:

-(13)

$$\underline{\sigma} \cdot \underline{\hat{p}} \phi_s^R = -i\hbar c \left(\frac{\partial}{\partial r} - \frac{j+1/2}{r} \right) F |jml_+ \rangle \quad -(14)$$

$$\underline{\sigma} \cdot \underline{\hat{p}} \phi_s^L = -\hbar c \left(\frac{\partial}{\partial r} + \frac{j+3/2}{r} \right) f |jml_- \rangle \quad -(15)$$

So eqs. (5) and (6) become:

$$(E - mc^2) F - \hbar c \left(\frac{d}{dr} + \frac{j+3/2}{r} \right) f = 0 \quad -(16)$$

$$(E + mc^2) f + \hbar c \left(\frac{d}{dr} - \frac{j-1/2}{r} \right) F = 0 \quad -(17)$$

which when solved simultaneously give the energy levels of H as in Merzbacher for example.

To derive eqs. (14) and (15) we:

$$\underline{\sigma} \cdot \underline{\hat{L}} |jml_- \rangle = (j - \frac{1}{2}) |jml_- \rangle \quad -(18)$$

$$\underline{\sigma} \cdot \underline{\hat{L}} |jml_+ \rangle = -(j + \frac{3}{2}) |jml_+ \rangle \quad -(19)$$

and note that:

$$\underline{\sigma} \cdot \underline{\hat{r}} \frac{\partial}{\partial r} = -\frac{\partial}{\partial r} \quad -(20)$$

$$\underline{\sigma} \cdot \underline{\hat{r}} \frac{1}{r} = -\frac{1}{r} \quad -(21)$$

because of parity.

Therefore:

$$\begin{aligned}
 3) \quad \underline{\sigma} \cdot \underline{\hat{p}} \phi_s^R &= \underline{\sigma} \cdot \underline{\hat{r}} \left(-i\hbar \frac{\partial}{\partial r} + i \frac{\underline{\sigma} \cdot \underline{\hat{L}}}{r} \right) F |j m l_+ \rangle \\
 &= -i\hbar \frac{\partial F}{\partial r} |j m l_+ \rangle + i\hbar \left(j - \frac{1}{2} \right) F |j m l_+ \rangle \\
 &= -i\hbar \left(\frac{\partial}{\partial r} - \frac{(j - 1/2)}{r} \right) F |j m l_+ \rangle
 \end{aligned}$$

$$\boxed{\underline{\sigma} \cdot \underline{\hat{p}} \phi_s^R = -i\hbar \left(\frac{\partial}{\partial r} - \frac{(j - 1/2)}{r} \right) F |j m l_+ \rangle}$$

and

$$\begin{aligned}
 \underline{\sigma} \cdot \underline{\hat{p}} \phi_s^L &= -i \underline{\sigma} \cdot \underline{\hat{p}} f |j m l_+ \rangle \\
 &= -i (\underline{\sigma} \cdot \underline{\hat{r}}) \left(-i\hbar \frac{\partial}{\partial r} + i \frac{\underline{\sigma} \cdot \underline{\hat{L}}}{r} \right) f |j m l_+ \rangle \\
 &= -\hbar \frac{\partial f}{\partial r} |j m l_- \rangle - \hbar \left(j + \frac{3}{2} \right) \frac{f}{r} |j m l_- \rangle \\
 &= -\hbar \left(\frac{\partial}{\partial r} + \frac{j + 3/2}{r} \right) f |j m l_- \rangle
 \end{aligned}$$

$$\boxed{\underline{\sigma} \cdot \underline{\hat{p}} \phi_s^L = -\hbar \left(\frac{\partial}{\partial r} + \frac{j + 3/2}{r} \right) f |j m l_- \rangle}$$

The complete wavefunctions ϕ^L and ϕ^R of the fermion equation are the combinations (3) and (4) in which ϕ^L and ϕ^R are interconverted by parity. Therefore:

4)

$$\phi^L = \phi_S^R + \phi_S^L$$

$$= F |j m l - \rangle - i f |j m l + \rangle \quad - (22)$$

$$\phi^R = \phi_S^R - \phi_S^L$$

$$= F |j m l - \rangle + i f |j m l + \rangle \quad - (23)$$

and $\hat{p} \phi^L = \phi^R \quad - (24)$

Therefore the effect of p on ϕ^L is to produce the complex conjugate.

Using De Moivre theorem:

$$e^{i\theta} = \cos \theta + i \sin \theta \quad - (25)$$

$$e^{-i\theta} = \cos \theta - i \sin \theta \quad - (26)$$

then:

$$\boxed{\phi^L = e^{i\theta}, \quad \phi^R = e^{-i\theta}} \quad - (27)$$

where $\cos \theta = F |j m l - \rangle \quad - (28)$

$$\sin \theta = f |j m l + \rangle \quad - (29)$$

So in eqs. (1) and (2):

$$(E + \underline{\sigma} \cdot \underline{p}) e^{i\theta} = m c^2 e^{-i\theta} \quad - (30)$$

$$(E - \underline{\sigma} \cdot \underline{p}) e^{-i\theta} = m c^2 e^{i\theta} \quad - (31)$$

and $\phi_S^R = \cos \theta, \quad \phi_S^L = -i \sin \theta \quad - (32)$