

186(7) : Calculation of Christoffel Symbols Directly
for Metric Compatibility.

The metric compatibility condition is:

$$\partial_\rho g_{\mu\nu} - \Gamma_{\rho\mu}^\lambda g_{\lambda\nu} - \Gamma_{\rho\nu}^\lambda g_{\mu\lambda} = 0 \quad - (1)$$

Consider a metric of the type:

$$g_{\mu\nu} = \begin{bmatrix} 1 - \frac{r_0}{r} & 0 & 0 & 0 \\ 0 & (1 - \frac{r_0}{r})^{-1} & 0 & 0 \\ 0 & 0 & -r^2 & 0 \\ 0 & 0 & 0 & -r^2 \sin^2 \phi \end{bmatrix} \quad - (2)$$

$$g^{\mu\nu} = \begin{bmatrix} (1 - \frac{r_0}{r})^{-1} & 0 & 0 & 0 \\ 0 & (1 - \frac{r_0}{r}) & 0 & 0 \\ 0 & 0 & -1/r^2 & 0 \\ 0 & 0 & 0 & -1/(r^2 \sin^2 \phi) \end{bmatrix} \quad - (3)$$

For $\mu = \nu = 0 \quad - (4)$

eq. (1) becomes:

$$\partial_\rho g_{00} - \Gamma_{\rho 0}^0 g_{00} - \Gamma_{\rho 0}^0 g_{00} \quad - (5)$$

i.e. $\Gamma_{10}^0 = \frac{1}{2g_{00}} \partial_1 g_{00} \quad - (6)$

$$= \frac{1}{2(1 - \frac{r_0}{r})} \frac{\partial}{\partial r} \left(1 - \frac{r_0}{r} \right)$$

$$\Gamma_{10}^0 = \frac{r_0}{2r^2 \left(1 - \frac{r_0}{r}\right)} \quad - (7)$$

for $\mu = \nu = 1$

$$\partial_\rho g_{11} - \Gamma_{\rho 1}^1 g_{11} - \Gamma_{\rho 1}^1 g_{11} = 0 \quad - (8)$$

i.e. $\partial_1 g_{11} = 2 \Gamma_{11}^1 g_{11} \quad - (9)$

for $\mu = \nu = 2 \quad - (10)$

$$\partial_1 g_{22} = 2 \Gamma_{12}^2 g_{22} \quad - (11)$$

$$\Gamma_{12}^2 = \frac{1}{2g_{22}} \partial_1 g_{22} \quad - (12)$$

$$\Gamma_{12}^2 = \frac{1}{r} \quad - (13)$$

For $\mu = \nu = 3 \quad - (14)$

$$\partial_\rho g_{33} = 2 \Gamma_{\rho 3}^3 g_{33} \quad - (15)$$

and $\Gamma_{13}^3 = \frac{1}{2g_{33}} \partial_1 g_{33}$

$$\Gamma_{13}^3 = \frac{1}{r} \quad - (16)$$

$$3) \quad \Gamma_{23}^3 = \frac{1}{2g_{33}} \partial_2 g_{33} = \frac{\sin \phi}{\cos \phi} \quad - (17)$$

In the limit: $r_0 \ll r \quad - (18)$

$$\Gamma_{10}^0 \rightarrow \frac{MG}{c^2 r^2} \quad - (19)$$

and

$$F = -mc^2 \Gamma_{10}^0 = -\frac{mMG}{r^2} \quad - (20)$$

From eq. (9) it is seen that Γ_{11}^1 exists by metric compatibility but the commutator method shows that for Γ_{11}^1 both the torsion and curvature vanish. So Γ_{11}^1 is a property of a spacetime in which there is no curvature or torsion.

In gravitational theory, only Γ_{10}^0 is significant, and for disconnection:

$$T_{10}^0 = 2\Gamma_{10}^0 \quad - (21)$$

and the curvature is zero.