

201 (1) : Some Transformation Properties of the New Relativity.

The covariant derivative is a derivative that is covariant under the general coordinate transformation:

$$D_{\mu} V^{\nu'} = \frac{\partial x^{\mu}}{\partial x^{\mu'}} \frac{\partial x^{\nu'}}{\partial x^{\nu}} D_{\mu} V^{\nu} \quad - (1)$$

This is a tensorial transformation law. It is defined by:

$$D_{\mu} V^{\nu} = \partial_{\mu} V^{\nu} + \Gamma^{\nu}_{\mu\lambda} V^{\lambda} \quad - (2)$$

where $\Gamma^{\nu}_{\mu\lambda}$ is the Christoffel connection. The covariant derivative is a tensor that reduces to the partial derivative of a vector in flat spacetime and transforms as a tensor in the arbitrary manifold.

In four dimensions $\Gamma^{\nu}_{\mu\lambda}$ has 64 components in general, and in general has no symmetry. In general $\Gamma^{\nu}_{\mu\lambda}$ has no symmetry in its lower two indices, i.e. asymmetric in its lower two indices μ and λ .

The ordinary derivative $\partial_{\mu} V^{\nu}$ does not transform as a tensor, because:

$$\begin{aligned} \frac{\partial}{\partial x^{\mu'}} (V^{\nu'}) &= \frac{\partial x^{\mu}}{\partial x^{\mu'}} \frac{\partial}{\partial x^{\mu}} \left(\frac{\partial x^{\nu'}}{\partial x^{\nu}} V^{\nu} \right) \quad - (3) \\ &= \frac{\partial x^{\mu}}{\partial x^{\mu'}} \frac{\partial}{\partial x^{\mu}} \left(\frac{\partial x^{\nu'}}{\partial x^{\nu}} \right) V^{\nu} + \frac{\partial x^{\mu}}{\partial x^{\mu'}} \frac{\partial x^{\nu'}}{\partial x^{\nu}} \frac{\partial V^{\nu}}{\partial x^{\mu}} \end{aligned}$$

however, for flat space:

$$2) \quad \frac{\partial}{\partial x^\mu} \left(\frac{\partial x^{\nu'}}{\partial x^\nu} \right) = 0 \quad - (4)$$

So in flat space:

$$\partial_\mu V^{\nu'} = \frac{\partial x^\mu}{\partial x^{\mu'}} \frac{\partial x^{\nu'}}{\partial x^\nu} \partial_\mu V^\nu \quad - (5)$$

and this is a tensor transformation.

The transformation law for the connection is derived from:

$$\begin{aligned} \partial_{\mu'} V^{\nu'} + \Gamma_{\mu'\lambda'}^{\nu'} V^{\lambda'} &= \frac{\partial x^\mu}{\partial x^{\mu'}} \frac{\partial x^{\nu'}}{\partial x^\nu} \partial_\mu V^\nu + \frac{\partial x^\mu}{\partial x^{\mu'}} V^\nu \frac{\partial}{\partial x^{\mu'}} \frac{\partial x^{\nu'}}{\partial x^\nu} + \Gamma_{\mu'\lambda'}^{\nu'} \frac{\partial x^{\lambda'}}{\partial x^\lambda} V^\lambda \\ &= \frac{\partial x^\mu}{\partial x^{\mu'}} \frac{\partial x^{\nu'}}{\partial x^\nu} \partial_\mu V^\nu \\ &= \frac{\partial x^\mu}{\partial x^{\mu'}} \frac{\partial x^{\nu'}}{\partial x^\nu} (\partial_\mu V^\nu + \Gamma_{\mu\lambda}^\nu V^\lambda) \end{aligned} \quad - (6)$$

So:

$$\frac{\partial x^\mu}{\partial x^{\mu'}} V^\nu \frac{\partial}{\partial x^{\mu'}} \frac{\partial x^{\nu'}}{\partial x^\nu} + \Gamma_{\mu'\lambda'}^{\nu'} \frac{\partial x^{\lambda'}}{\partial x^\lambda} V^\lambda = \frac{\partial x^\mu}{\partial x^{\mu'}} \frac{\partial x^{\nu'}}{\partial x^\nu} \Gamma_{\mu\lambda}^\nu V^\lambda \quad - (7)$$

In the first term on the LHS change summation index from ν to λ :

$$\frac{\partial x^\mu}{\partial x^{\mu'}} V^\lambda \frac{\partial}{\partial x^{\mu'}} \frac{\partial x^{\nu'}}{\partial x^\lambda} = \frac{\partial x^\mu}{\partial x^{\mu'}} V^\nu \frac{\partial}{\partial x^{\mu'}} \frac{\partial x^{\nu'}}{\partial x^\nu} \quad - (8)$$

Then V^λ cancels from eq. (7) to give:

$$\frac{\partial x^\mu}{\partial x^{\mu'}} \frac{\partial}{\partial x^\mu} \frac{\partial x^{\nu'}}{\partial x^\lambda} + \Gamma_{\mu'\lambda'}^{\nu'} \frac{\partial x^{\lambda'}}{\partial x^\lambda} = \frac{\partial x^\mu}{\partial x^{\mu'}} \frac{\partial x^{\nu'}}{\partial x^\nu} \Gamma_{\mu\lambda}^\nu - (9)$$

Finally multiply by $\partial x^{\lambda'}/\partial x^{\lambda'}$ to give

$$\Gamma_{\mu'\lambda'}^{\nu'} = \frac{\partial x^\mu}{\partial x^{\mu'}} \frac{\partial x^{\lambda'}}{\partial x^{\lambda'}} \frac{\partial x^{\nu'}}{\partial x^\nu} \Gamma_{\mu\lambda}^\nu - \frac{\partial x^\mu}{\partial x^{\mu'}} \frac{\partial x^{\lambda'}}{\partial x^{\lambda'}} \frac{\partial}{\partial x^\mu} \left(\frac{\partial x^{\nu'}}{\partial x^{\lambda'}} \right)$$

As in the appended notes it follows that the difference of any two connections is a tensor:

$$S_{\mu\nu}^\lambda = \Gamma_{\mu\nu}^\lambda - \hat{\Gamma}_{\mu\nu}^\lambda - (11)$$

In general relativity, the quantities of physics must be tensors. So the second, inhomogeneous, term in eq. (10) is unphysical.

Its symmetry is:

$$\frac{\partial x^\mu}{\partial x^{\mu'}} \frac{\partial x^{\lambda'}}{\partial x^{\lambda'}} \frac{\partial^2 x^{\nu'}}{\partial x^\mu \partial x^{\lambda'}} = \frac{\partial x^{\lambda'}}{\partial x^{\lambda'}} \frac{\partial x^\mu}{\partial x^{\mu'}} \frac{\partial^2 x^{\nu'}}{\partial x^{\lambda'} \partial x^\mu} - (12)$$

but this symmetry is irrelevant to physics.

It is well known that:

$$[D_\mu, D_\nu] V^\rho = R^\rho_{\sigma\mu\nu} V^\sigma - (\Gamma_{\mu\nu}^\lambda - \Gamma_{\nu\mu}^\lambda) D_\lambda V^\rho$$

definition:

$$[D_\mu, D_\nu] = -[D_\nu, D_\mu] - (13)$$

$$[D_\mu, D_\nu] = -[D_\nu, D_\mu] - (14)$$

4) $R^{\rho}_{\sigma\mu\nu} = -R^{\rho}_{\sigma\nu\mu}$ — (15)
 by definition. The commutator and curvature represent a
 round trip in the abstract manifold.
 Therefore, by definition, the connection is anti-
symmetric: $\Gamma^{\lambda}_{\mu\nu} = -\Gamma^{\lambda}_{\nu\mu}$ — (16)

When: $\mu = \nu$ — (17)

$$[D_{\mu}, D_{\nu}] = [D_{\mu}, D_{\nu}] = 0 \quad - (18)$$

and $\Gamma^{\lambda}_{\mu\nu} (\mu = \nu) = 0$ — (19)

Similarly, $R^{\rho}_{\sigma\mu\nu} (\mu = \nu) = 0$ — (20)

and $T^{\lambda}_{\mu\nu} = \Gamma^{\lambda}_{\mu\nu} - \Gamma^{\lambda}_{\nu\mu} = -T^{\lambda}_{\nu\mu}$ — (21)

$T^{\lambda}_{\mu\nu} (\mu = \nu) = 0$ — (22)

The connection is antisymmetric in all frames ✓
 reference. It follows that:

$$\frac{\partial x^{\mu}}{\partial x^{\mu'}} \frac{\partial x^{\lambda}}{\partial x^{\lambda'}} \frac{\partial^2 x^{\nu'}}{\partial x^{\mu} \partial x^{\lambda}} = 0 \quad - (23)$$

because it is symmetric, and the symmetric part
 of the connection is zero.

3) The antisymmetric connection is a tensor:

$$\boxed{\Gamma^{\nu'}_{\mu'\lambda'} = \frac{\partial x^\mu}{\partial x^{\mu'}} \frac{\partial x^\lambda}{\partial x^{\lambda'}} \frac{\partial x^{\nu'}}{\partial x^\nu} \Gamma^{\nu}_{\mu\lambda}} \quad - (4)$$

and the antisymmetric connection is physical.

The transformation properties of the equivalence theorem of UFT 199 now follow. The equivalence theorem is:

$$d_\mu \nabla^a = \omega^a_{\mu\nu} - (5)$$

and

$$d_\mu \nabla^a = \omega^a_{\mu b} \nabla^b - (6)$$

where

$$D_\mu \nabla^a = d_\mu \nabla^a + \omega^a_{\mu b} \nabla^b - (7)$$

Eq. (6) implies:

$$\boxed{d_\mu \nabla^\nu = \Gamma^{\nu}_{\mu\lambda} \nabla^\lambda} \quad - (8)$$

static frame,
vector analysis,
flat space

static vector,
general manifold,
ordinary derivative zero

Both sides of eq. (8) are tensorial because:

$$d_{\mu'} \nabla^{\nu'} = \frac{\partial x^\mu}{\partial x^{\mu'}} \frac{\partial x^{\nu'}}{\partial x^\nu} d_\mu \nabla^\nu - (9)$$

$$\Gamma^{\nu'}_{\mu'\lambda'} \nabla^{\lambda'} = \Gamma^{\nu'}_{\mu'\lambda'} \frac{\partial x^{\lambda'}}{\partial x^\lambda} \nabla^\lambda - (10)$$

Transformation of the Connection

$$\begin{aligned} \Gamma_{\mu'\lambda'}^{\nu'} &= \frac{\partial x^\mu}{\partial x^{\mu'}} \frac{\partial x^\lambda}{\partial x^{\lambda'}} \frac{\partial x^{\nu'}}{\partial x^\nu} \Gamma_{\mu\lambda}^\nu \\ &\quad - \frac{\partial x^\mu}{\partial x^{\mu'}} \frac{\partial x^\lambda}{\partial x^{\lambda'}} \frac{\partial^2 x^{\nu'}}{\partial x^\mu \partial x^\lambda} \end{aligned} \quad (1)$$

Symmetries

In general, $\Gamma_{\mu'\lambda'}^{\nu'}$ is asymmetric in its lower two indices, i.e. in isolation of any other argument it has no symmetry in μ' and λ' or μ and λ . So the tensor is :

$$\begin{aligned} T_{\mu'\lambda'}^{\nu'} &= \Gamma_{\mu'\lambda'}^{\nu'} - \Gamma_{\lambda'\mu'}^{\nu'} \\ &= \frac{\partial x^\mu}{\partial x^{\mu'}} \frac{\partial x^\lambda}{\partial x^{\lambda'}} \frac{\partial x^{\nu'}}{\partial x^\nu} \left(\Gamma_{\mu\lambda}^\nu - \Gamma_{\lambda\mu}^\nu \right) \\ &= \left(\frac{\partial x^\mu}{\partial x^{\mu'}} \frac{\partial x^\lambda}{\partial x^{\lambda'}} \frac{\partial x^{\nu'}}{\partial x^\nu} \right) T_{\mu\lambda}^\nu \end{aligned} \quad (2)$$

and is a tensor which is generally covariant. This is very well known.

For any symmetry :

$$\frac{\partial x^\mu}{\partial x^{\mu'}} \frac{\partial x^\lambda}{\partial x^{\lambda'}} \frac{\partial^2 x^{\nu'}}{\partial x^\mu \partial x^\lambda} = \frac{\partial x^\lambda}{\partial x^{\lambda'}} \frac{\partial x^\mu}{\partial x^{\mu'}} \frac{\partial^2 x^{\nu'}}{\partial x^\lambda \partial x^\mu} \quad (3)$$

is very well known.

THE ARGUMENTS OF RODRIGUES ARE INCORRECT

The details of eq. (3) are as follows:

$$1) \quad \frac{\partial x^\lambda}{\partial x^{\lambda'}} \frac{\partial x^\mu}{\partial x^{\mu'}} = \frac{\partial x^\mu}{\partial x^{\mu'}} \frac{\partial x^\lambda}{\partial x^{\lambda'}} \quad - (4)$$

This is an algebraic result:

$$ab = ba \quad - (5)$$

$$2) \quad \frac{\partial^2 x^{\nu'}}{\partial x^\lambda \partial x^\mu} = \frac{\partial^2 x^{\nu'}}{\partial x^\mu \partial x^\lambda} \quad - (6)$$

is a property of the ordinary derivative.

In general, a large number of conventions can be defined on a manifold. For any two conventions, their difference is a (1, 2) tensor (S.P. Carroll "Spacetime and Geometry: an Introduction to General Relativity" (Addison Wesley, NY, 2004). So it is general:

$$S_{\mu\nu}^\lambda = \Gamma_{\mu\nu}^\lambda - \hat{\Gamma}_{\mu\nu}^\lambda \quad - (7)$$

is a tensor which transforms as:

$$\begin{aligned} S_{\mu'\nu'}^{\lambda'} &= \Gamma_{\mu'\nu'}^{\lambda'} - \hat{\Gamma}_{\mu'\nu'}^{\lambda'} \\ &= \frac{\partial x^\mu}{\partial x^{\mu'}} \frac{\partial x^\nu}{\partial x^{\nu'}} \frac{\partial x^{\lambda'}}{\partial x^\lambda} \Gamma_{\mu\nu}^\lambda - \frac{\partial x^\mu}{\partial x^{\mu'}} \frac{\partial x^\nu}{\partial x^{\nu'}} \frac{\partial^2 x^{\lambda'}}{\partial x^\mu \partial x^\nu} \\ &\quad - \frac{\partial x^\mu}{\partial x^{\mu'}} \frac{\partial x^\nu}{\partial x^{\nu'}} \frac{\partial x^{\lambda'}}{\partial x^\lambda} \hat{\Gamma}_{\mu\nu}^\lambda + \frac{\partial x^\mu}{\partial x^{\mu'}} \frac{\partial x^\nu}{\partial x^{\nu'}} \frac{\partial^2 x^{\lambda'}}{\partial x^\mu \partial x^\nu} \\ &= \frac{\partial x^\mu}{\partial x^{\mu'}} \frac{\partial x^\nu}{\partial x^{\nu'}} \frac{\partial x^{\lambda'}}{\partial x^\lambda} (\Gamma_{\mu\nu}^\lambda - \hat{\Gamma}_{\mu\nu}^\lambda) \\ &= \frac{\partial x^\mu}{\partial x^{\mu'}} \frac{\partial x^\nu}{\partial x^{\nu'}} \frac{\partial x^{\lambda'}}{\partial x^\lambda} S_{\mu\nu}^\lambda \quad - (8) \end{aligned}$$

3) Levi-Civita transformation law (Q.E.D.)

Note carefully that the difference of any two connections is a tensor. This result does not depend on the symmetry of the lower two indices of the connection.

Now consider the special case:

$$\hat{\Gamma}_{\mu\nu}^{\lambda} = \Gamma_{\nu\mu}^{\lambda} \quad - (9)$$

It follows that:

$$T_{\mu\nu}^{\lambda} = \Gamma_{\mu\nu}^{\lambda} - \Gamma_{\nu\mu}^{\lambda} \quad - (10)$$

is a tensor. It is known as the torsion tensor. From consideration of the commutator.

$$[D_{\mu}, D_{\nu}] V^{\rho} = R^{\rho}_{\sigma\mu\nu} V^{\sigma} - T_{\mu\nu}^{\lambda} D_{\lambda} V^{\rho} \quad - (11)$$

where $[D_{\mu}, D_{\nu}] = -[D_{\nu}, D_{\mu}] \quad - (12)$

Eq. (11) is:

$$[D_{\mu}, D_{\nu}] V^{\rho} = -(\Gamma_{\mu\nu}^{\lambda} - \Gamma_{\nu\mu}^{\lambda}) D_{\lambda} V^{\rho} + R^{\rho}_{\sigma\mu\nu} V^{\sigma} \quad - (13)$$

As is well known: $R^{\rho}_{\sigma\mu\nu} = -R^{\rho}_{\sigma\nu\mu} \quad - (14)$

so

$$\Gamma_{\mu\nu}^{\lambda} = -\Gamma_{\nu\mu}^{\lambda} \quad - (15)$$