

242(7) : Baseline Check by Computer

From note 242(6) the true anomaly is:

$$\theta = \sqrt{2} \frac{A}{T} \int \frac{f(r)}{r^2} dr \quad - (1)$$

where $f(r) = \left(- \int r \Omega^2 dr - \infty \right)^{-1/2} \quad - (2)$

and $\Omega^2 = - \frac{L_0^2}{m^2 r^4} - \frac{F(r)}{mr} \quad - (3)$

The Sacc wave equation is:

$$\frac{d^2 r}{dt^2} + \Omega^2 r = 0 \quad - (4)$$

The Newtonian force is:

$$F(r) = - \frac{L_0^2}{dmr^2} \quad - (5)$$

in which case:

$$f(r) = \left(\frac{2L_0^2 r - L_0^2 d}{2dm^2 r^2} - A \right)^{-1/2} \quad - (6)$$

For the square root to be real:

$$L_0^2 > 2d^2 \times m^2 \quad - (7)$$

The Newtonian result from eq. (1) is

therefore:

$$\theta = \gamma - \frac{2m A}{L_0 T} \sin^{-1} \left(\frac{\frac{d}{r} - 1}{(1 - 2d^2 \alpha m^2)^{1/2}} \right) \quad - (8)$$

If:

$$x = - \frac{L_0^2}{2d^2 m^2} (\epsilon - 1)(\epsilon + 1), \quad - (9)$$

$$\gamma = \pi \quad - (10)$$

Then:

$$\theta = \frac{m A}{L_0 T} \left(\pi - 2 \sin^{-1} \left(\frac{\frac{d}{r} - 1}{((\epsilon - 1)(\epsilon + 1) + 1)^{1/2}} \right) \right) \quad - (11)$$

i.e.:

$$r = \frac{d}{1 + \epsilon \cos \left(\frac{L_0 T}{2m A} \theta \right)} \quad - (12)$$

For any curve in Lagrangian dynamics:

$$dt = \frac{2m}{L_0} dA \quad - (13)$$

3) (J. D. Marra and S. T. Thornton, "Classical Dynamics" (Harvard, 1988, 3rd ed.))

The area of the ellipse is swept out in time T , so

$$\int_0^T dt = \frac{2m}{L_0} \int_0^A dA \quad - (14)$$

so

$$T = \frac{2m}{L_0} A \quad - (15)$$

From eqs. (12) and (15):

$$r = \frac{d}{1 + \epsilon \cos \theta} \quad - (16)$$

QED

In Newtonian dynamics, the orbit is an

ellipse with area:

$$A = \pi ab \quad - (17)$$

where

$$a = \frac{d}{1 - \epsilon^2} = \frac{n M G}{2 |E|} \quad - (18)$$

$$b = \frac{d}{(1 - \epsilon^2)^{1/2}} = \frac{L_0}{(2m |E|)^{1/2}} \quad - (19)$$

So:

$$T = \frac{2m}{L_0} \cdot \pi ab = \pi n M G \left(\frac{n}{2} \right)^{1/2} E^{-3/2} \quad - (20)$$

4) The semimajor axis can be written as:

$$b = (d \cdot a)^{1/2} \quad - (21)$$

where

$$d = \frac{L_0^2}{m^2 \underline{M} G} \quad - (22)$$

so:

$$T^2 = \left(\frac{4\pi^2}{\underline{M} G} \right) a^3 \quad - (23)$$

which is Kepler's Third Law of Planetary Motion.

Eq (9) can be simplified by noting that:

$$a = \frac{d}{1-e^2} = \frac{d}{(1-e)(1+e)} \quad - (24)$$

so

$$x = \frac{L_0^2}{2adm^2} \quad - (25)$$

where:

$$L_0^2 = m^2 \underline{M} G d \quad - (26)$$

so

$$\boxed{x = \frac{\underline{M} G}{2a}} \quad - (27)$$

which has the correct units of the square of

5) velocity.

The true anomaly may now be calculated for any force law and any planar orbit. For example if:

$$F(r) = -\frac{mMGx_1^2}{r^2} - d(1-x_1^2)\frac{mMG}{r^3} \quad (28)$$

The orbit from Lagrangian dynamics is:

$$r = \frac{d}{1 + \epsilon \cos(x_1 \theta)} \quad (29)$$

where x_1 is the precession factor. By using eq. (28) in eq. (1), eq. (29) should result, and x_1 can be calculated exactly.

Einsteinian general relativity claims

that:

$$F(r) = -\frac{mMG}{r^2} - \frac{3L_0^2 mMG}{m^2 c^2 r^4} \quad (30)$$

The force law eq. (30) can now be used in eq. (1) to find the orbit, r as a function of θ .

6) The claim of EGR is that eq. (30) produces:

$$\theta = 2\pi (1 + x_1) \quad - (31)$$

per revolution of 2π , where:

$$x_1 = \frac{3GM}{ac^2(1-e^2)} \quad - (32)$$

from eqs. (15) and (11) the Newtonian result is:

$$\theta = \frac{1}{2} \left(\pi - 2 \sin^{-1} \left(\frac{\frac{d}{r} - 1}{1((e-1)(e+1)+1)^{1/2}} \right) \right) \quad - (33)$$

If

$$\theta = 2\pi \quad - (34)$$

the Newtonian result from eq. (16) is:

$$r = \frac{d}{1+e} \quad - (35)$$

According to EGR, the force law (30) must

produce: $\theta \rightarrow 2\pi \left(1 + \frac{3GM}{ac^2(1-e^2)} \right) \quad - (36)$

so using the force law (30) in eq. (1) must change eq. (35) to:

7)

$$r = \frac{d}{1 + \epsilon \cos(x, \theta)} \quad (37)$$

If we consider the general equation (1), a change is made from eq. (5) to eq. (30) should produce the result (37) if EGR is true. This is therefore a test of EGR.