

243(6): Some Implications of the Concept of Mean Curvature in Radiation Theory

The mean curvature in radiation theory was defined in the previous note as:

$$\langle R \rangle = \frac{1}{h^2 c^3} \langle E^2 \rangle \quad - (1)$$

where

$$\langle E^2 \rangle = \frac{h^2 \omega^2}{\exp\left(\frac{h\omega}{kT}\right) - 1} \quad - (2)$$

so

$$\langle R \rangle = \frac{\omega^2}{c^3} \left(\exp\left(\frac{h\omega}{kT}\right) - 1 \right)^{-1} \quad - (3)$$

The mean energy in the Planck theory is:

$$\langle E \rangle = \frac{h\omega}{\exp\left(\frac{h\omega}{kT}\right) - 1} \quad - (4)$$

so:

$$\langle E \rangle = \frac{h c^3}{\omega} \langle R \rangle \quad - (5)$$

From eqns. (1) and (5):

$$\langle E^2 \rangle = h\omega \langle E \rangle \quad - (6)$$

From the ECE wave equation:

2)

$$(\square + R) \gamma_\mu^a = 0 \quad - (7)$$

it may be inferred that:

$$R = \left(\frac{mc}{\hbar} \right)^2 \quad - (8)$$

So

$$\langle R \rangle = \langle m^2 \rangle \frac{c^2}{\hbar^2} \quad - (9)$$

From eqs. (1) and (9):

$$\langle m^2 \rangle = \frac{\hbar \omega}{c^4} \langle E \rangle \quad - (10)$$

In the Planck-Einstein theory $\langle m^2 \rangle$ is
the mean square photon mass.

For a photon at rest:

$$\omega = \omega_0 \quad - (11)$$

and

$$\langle E \rangle = \hbar \omega_0 \quad - (12)$$

For moving photons:

$$\langle m^2 \rangle = \frac{\hbar^2 \omega^2}{c^4 \left(\exp \left(\frac{\hbar \omega}{kT} \right) - 1 \right)} \quad - (13)$$

3) This is a new expression for mean square photon mass. In the older de Broglie / Einstein theory:

$$E = \gamma mc^2 = \hbar \omega \quad (14)$$

$$\underline{p} = \gamma m \underline{v} = \hbar \underline{k} \quad (15)$$

but as shown in HFT 158 ff., eqs. (14) and (15) fail completely in Compton scattering and related effects.

The density of states in the Planck Einstein theory is defined by:

$$\rho(\omega) = \frac{dU}{d\omega} = \frac{\hbar}{\pi^2 c^3} \frac{\omega^3}{\left(\exp\left(\frac{\hbar\omega}{kT}\right) - 1\right)}$$

photon

so mean square / mass is related to $\rho(\omega)$ density of states as follows:

$$\boxed{\omega c \pi^2 \langle m^2 \rangle = \hbar \rho(\omega)} \quad (17)$$

so the mean square photon mass is quantized in terms of the density of states

from eqs. (16) and (17):

4)

$$\frac{dU}{d\omega} = \frac{\pi^2 \omega c}{\hbar} \langle m^2 \rangle \quad - (18)$$

$$= \frac{\hbar}{\pi^2 c^3} \frac{\omega^3}{e^{\hbar\omega/kT} - 1} \quad - (19)$$

The total energy density of black body radiation is therefore:

$$\begin{aligned} U &= \frac{\hbar}{\pi^2 c^3} \int_0^\infty \frac{\omega^3}{e^{\hbar\omega/kT} - 1} d\omega \\ &= \frac{\pi^2 c}{\hbar} \int_0^\infty \omega \langle m^2 \rangle d\omega \quad - (20) \\ &= \left(\left(\frac{\pi^2}{15} \right) \frac{\hbar^4}{c^3 \hbar^3} \right) T^4 \end{aligned}$$

which relates to mean square photon mass to the Stefan Boltzmann law as follows:

$$\boxed{\int_0^\infty \omega \langle m^2 \rangle d\omega = \left(\frac{\hbar^4}{15 c^4 \hbar^3} \right) T^4} \quad - (21)$$

This shows that the Stefan Boltzmann law

5) is due to near square photon mass.

The mean energy of the black distribution is

$$\langle E \rangle = \frac{c^4}{\hbar \omega} \langle m^2 \rangle \quad - (22)$$

so the heat capacity is:

$$\begin{aligned} C_v &= 3N d \frac{\langle E \rangle}{dT} = \frac{3N c^4}{\hbar \omega} \frac{d \langle m^2 \rangle}{dT} \\ &= 3N k \left(\frac{\hbar \omega}{kT} \right)^2 \frac{e^{\hbar \omega / kT}}{(1 - e^{\hbar \omega / kT})^2} \quad - (23) \end{aligned}$$

so

$$\boxed{\frac{d \langle m^2 \rangle}{dT} = \frac{(\hbar \omega)^3}{(kT)^2} \frac{k}{c^4} \frac{e^{\hbar \omega / kT}}{(1 - e^{\hbar \omega / kT})^2}}$$

- (24)

which is the capacity of near square photon mass.
