

257(4): General Solution of the Beltrami Equation in
Terms of Elliptically Polarized Plane Waves.

The general solution can be written in terms of:

$$B_x = B_2 \sin \phi_z + B_3 \cos \phi_y \quad - (1)$$

$$B_y = B_3 \sin \phi_x + B_1 \cos \phi_z \quad - (2)$$

$$B_z = B_1 \sin \phi_y + B_2 \cos \phi_x \quad - (3)$$

so
$$\underline{B} = B_x \underline{i} + B_y \underline{j} + B_z \underline{k} \quad - (4)$$

and
$$\nabla \times \underline{B} = \lambda \underline{B} \quad - (5)$$

where λ is a constant. Eqs. (1) to (3) have the structure of ABC flow:

$$u_1 = A \sin z + C \cos y \quad - (6)$$

$$u_2 = B \sin x + A \cos z \quad - (7)$$

$$u_3 = C \sin y + B \cos x \quad - (8)$$

These are examples of 3D chaotic phenomena and
lagrangian turbulence.

Now use:

$$\cos \phi = \frac{1}{2} (e^{i\phi} + e^{-i\phi}) \quad - (9)$$

$$\sin \phi = -\frac{i}{2} (e^{i\phi} - e^{-i\phi}) \quad - (10)$$

and eq. (4) can be written as a sum of two terms as follows:

$$\begin{aligned}
\underline{B} &= B_2 \sin \phi_z \underline{i} + B_1 \cos \phi_z \underline{j} \\
&\quad + B_3 \cos \phi_y \underline{i} + B_1 \sin \phi_y \underline{k} \\
&\quad + B_3 \sin \phi_x \underline{j} + B_2 \cos \phi_x \underline{k} \quad \text{--- (11)} \\
&= \frac{1}{2} \left[(i B_2 + B_1 j) e^{-i\phi_z} + (-i B_2 + B_1 j) e^{i\phi_z} \right. \\
&\quad + (B_3 i - i B_1 k) e^{i\phi_y} + (B_3 i + i B_1 k) e^{-i\phi_y} \\
&\quad \left. + (B_2 k - i B_3 j) e^{i\phi_x} + (B_2 k + i B_3 j) e^{-i\phi_x} \right]
\end{aligned}$$

If the phases are defined as:

$$\phi_x = \omega_x t - k_x X \quad \text{--- (12)}$$

$$\phi_y = \omega_y t - k_y Y \quad \text{--- (13)}$$

$$\phi_z = \omega_z t - k_z Z \quad \text{--- (14)}$$

Eq. (11) represents a sum of three elliptically polarized waves in X , Y and Z .

Each of the six terms in eq. (11) obey a Helmholtz equation and are also solutions to the vacuum equations of the simplified ECE engineering model.

There are also several conjugate products possible among the terms

3) For example, if:

$$\underline{B}^{(1)} = (iB_2 \underline{i} + B_1 \underline{j}) e^{-i\phi z} = \underline{B}^{(2)*} \quad - (15)$$

then:

$$\underline{B}^{(1)} \times \underline{B}^{(2)} = \begin{vmatrix} \underline{i} & \underline{j} & \underline{k} \\ iB_2 & B_1 & 0 \\ -iB_2 & B_1 & 0 \end{vmatrix} = 2iB_1B_2 \underline{k}, \quad - (16)$$

giving

$$\underline{B}_z^{(3)} = 2B_1B_2 \underline{k} \quad - (17)$$

Similarly

$$\underline{B}_y^{(3)} = 2B_1B_3 \underline{j} \quad - (18)$$

and

$$\underline{B}_x^{(3)} = 2B_2B_3 \underline{i} \quad - (19)$$

Since these equations have the structure of eqns. (6) to (8), i.e. of ABC flows, a transition to turbulence should occur under suitably defined conditions

Finally it is possible to consider a plane wave of the type:

$$\underline{B}^{(1)} = \frac{B(z)}{\sqrt{2}} (\underline{i} - i\underline{j}) e^{i(\omega t - k z)} \quad - (20)$$

but this

gives a type of non-Beltrami equation.

$$4) \quad \underline{\nabla} \times \underline{B}^{(1)} = \frac{1}{\sqrt{2}} \begin{vmatrix} \underline{i} & \underline{j} & \underline{k} \\ \partial/\partial x & \partial/\partial y & \partial/\partial z \\ B^{(1)}(z)e^{i\phi} & -iB^{(1)}(z)e^{i\phi} & 0 \end{vmatrix} \quad - (21)$$

$$= \kappa \underline{B}^{(1)} + \frac{1}{\sqrt{2}} \frac{\partial B(z)}{\partial z} (\underline{i} + i\underline{j}) e^{i\phi}$$

i.e.

$$\underline{\nabla} \times \underline{B}_L^{(1)} = \kappa \underline{B}_L^{(1)} + \frac{\partial B(z)}{\partial z} \underline{B}_R^{(1)} \quad - (22)$$

where:

$$\underline{B}_L^{(1)} = \frac{1}{\sqrt{2}} (\underline{i} - i\underline{j}) e^{i\phi} \quad - (23)$$

$$\underline{B}_R^{(1)} = \frac{1}{\sqrt{2}} (\underline{i} + i\underline{j}) e^{i\phi} \quad - (24)$$

Therefore the coefficients B_1 , B_2 and B_3 must be independent of x , y and z .