

## 263(1): A New Theory of Orbital Dynamics

In UFT 262 it has been shown that the experimentally observed precession of orbits is given by the equation:

$$d - r_0 = \frac{d}{1 + \epsilon \cos \theta} \quad - (1)$$

where

$$r_0 = \frac{3MG}{c^2} \quad - (2)$$

For a given  $M$ ,  $r_0$  is a universal constant. Here  $M$  is the mass at the focus of the ellipse,  $d$  is the half right latus recte,  $\epsilon$  the eccentricity and  $(r, \theta)$  the plane polar coordinate system. Eq. (1) was derived from the turning point of the obsolete Einstein theory:

$$\frac{d^2 r}{dt^2} = 0 = -m \frac{MG}{r^2} - \frac{3MG L^2}{m c^2 r^4} + \frac{L^2}{m r^3} \quad - (3)$$

giving:

$$\Delta \theta = \frac{3MG}{\epsilon d c^2} = \frac{3MG}{a c^2 (1 - \epsilon^2)} \quad - (4)$$

where  $a$  is the semi-major axis of the ellipse.

Unfortunately the Einstein theory is flawed due to its use of an obsolete geometry without revision.

There is, however, a much simpler and more direct method of deriving eq. (1). This

a) method can be developed from the definition:

$$\underline{R} = R \underline{e}_r \quad - (5)$$

where

$$R = r + r_0 \quad - (6)$$

Therefore in fundamental kinematic theory,  $r$  is replaced wherever it occurs by  $R$ . The coordinate system becomes  $(R, \theta)$ . Therefore the linear velocity is:

$$\underline{v} = \frac{dR}{dt} \underline{e}_r + R \frac{d\underline{e}_r}{dt} = \frac{dR}{dt} \underline{e}_r + \dot{\theta} R \underline{e}_\theta \quad - (7)$$

and the linear acceleration is:

$$\underline{a} = (\ddot{R} - R\dot{\theta}^2) \underline{e}_r + (R\ddot{\theta} + 2\dot{R}\dot{\theta}) \underline{e}_\theta \quad - (8)$$

In all planar orbits:

$$R\ddot{\theta} + 2\dot{R}\dot{\theta} = 0 \quad - (9)$$

So:

$$\underline{F}(R) = (\ddot{R} - R\dot{\theta}^2) \underline{e}_r \quad - (10)$$

$$= - \frac{L^2}{mR^2} \left( \frac{d^2}{d\theta^2} \left( \frac{1}{R} \right) + \frac{1}{R} \right) \underline{e}_r$$

and

$$\frac{d^2 R}{dt^2} = F(R) + mR\dot{\theta}^2 \quad - (11)$$

Eqs. (5) to (11) are true for any force law.

3) The turning point of the orbit is :

$$\frac{d^2 R}{dt^2} = 0 \quad - (12)$$

i.e.  $F(R) + m R \dot{\theta}^2 = 0 \quad - (13)$

The angular momentum is defined by :

$$L = m R^2 \dot{\theta} = \text{constant} \quad - (14)$$

where  $\omega = \dot{\theta} = \frac{d\theta}{dt} = \frac{L}{m R^2} \quad - (15)$

is the angular velocity and spin connection.

Now use the Hooke / Newton inverse square

law modified for  $R$  :

$$\underline{F} = - \frac{m M G}{R^2} \underline{e}_r \quad - (16)$$

Eq. (16) gives the modified elliptical orbit :

$$R = \frac{d}{1 + \epsilon \cos \theta} \quad - (17)$$

The turning point of this orbit is defined by :

$$- \frac{m M G}{R^2} + \frac{L^2}{m R^3} = 0 \quad - (18)$$

i.e.  $R = \frac{L^2}{m^2 M G} \quad - (19)$

4) From eq. (10) :

$$R = d = \frac{L^2}{m^2 M G} \quad - (20)$$

So at the turning point:

$$r + r_0 = d \quad - (21)$$

so 
$$r = d - r_0 \quad - (22)$$

Eq. (22) implies eq. (1), so the precession is the experimentally observed :

$$\Delta \theta = \frac{3 M G}{c^2 d} \quad - (23)$$

Eq. (23) has been derived entirely without the use of the Einstein theory

## Discussion

The above equation are examples of ECE theory with spin correction magnitude  $\omega$ . The experimental precession (23) is given by the spin correction :

$$\omega = \frac{L}{m R^2} = \frac{L}{m (r + r_0)^2} \quad - (24)$$

The Hooke / Newton / Kepler ellipse is

→ given by the spin correction:

$$\omega = \frac{L}{mr^2} \quad - (25)$$

The Hooke/Newton law is modified to:

$$F = -\frac{mM\bar{G}}{(r+r_0)^2} \sim -\frac{mM\bar{G}}{r^2} + \frac{2mM\bar{G}r_0}{r^3} \quad - (26)$$

which gives the precessing ellipse:

$$r = \frac{d}{1 + \epsilon \cos(x\theta)} \quad - (27)$$

where

$$x = 1 + \frac{r_0}{d} \quad - (28)$$

Using these equations the entire Einstein theory becomes irrelevant, and the entire subject of classical dynamics can be redeveloped with  $R$  replacing  $r$ .

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