

268(2) : Derivation of the Ellipse from the Classical Hamiltonian and Result from Quantum Theory

The classical Hamiltonian is :

$$H = E = \frac{p^2}{2m} - \frac{k}{r} \quad - (1)$$

as in previous notation. It can be expressed as :

$$E = \frac{m}{2} \left(\left(\frac{dr}{dt} \right)^2 + r^2 \left(\frac{d\theta}{dt} \right)^2 \right) - \frac{k}{r} \quad - (2)$$

so

$$\left(\frac{dr}{d\theta} \right)^2 = \frac{2mr^4}{L^2} \left(E + \frac{k}{r} \right) - r^2 \quad - (3)$$

using :

$$\frac{d\theta}{dt} = \frac{L}{mr^2} \quad - (4)$$

and

$$\frac{dr}{dt} = \frac{dr}{d\theta} \frac{d\theta}{dt} \quad - (5)$$

Now consider the ellipse :

$$r = \frac{a}{1 + e \cos \theta} \quad - (6)$$

It follows that :

$$\left(\frac{dr}{d\theta} \right)^2 = \frac{e^2 r^4}{a^2} (1 - \cos^2 \theta) \quad - (7)$$

where

$$\cos^2 \theta = \frac{1}{e^2} \left(\frac{a}{r} - 1 \right)^2 \quad - (8)$$

After some algebra :

$$2) \left(\frac{dr}{dt} \right)^2 = \frac{r^2}{L^2} \left(r^2 (\epsilon^2 - 1) + 2dr \right) - r^2 \quad - (9)$$

Comparing Eqs (3) and (9):

$$\boxed{\alpha = \frac{L^2}{mk}} \quad - (10)$$

and

$$\frac{\epsilon^2 - 1}{L^2} = \frac{2mE}{L^2} \quad - (11)$$

so

$$\epsilon^2 = 1 + \frac{2mEL^2}{L^2}$$

$$\boxed{\epsilon^2 = 1 + \frac{2EL^2}{mk^2}} \quad - (12)$$

It is seen that the ellipse (6) is a direct re-expression of the fundamental classical Hamiltonian (1).

This is true S.O. for gravitation and electrodynamics. For gravitation:

$$k = mM\gamma \quad - (13)$$

and for electrodynamics:

$$k = \frac{e}{4\pi\epsilon_0} \quad - (14)$$

3) Note that the Bohr theory of the atom is obtained directly by using:

$$L = n\hbar \quad - (15)$$

and

$$E = 0. \quad - (16)$$

In the Schrodinger theory the following results have been computed in recent work:

$$1) \quad \langle r \rangle = \left\langle \frac{a}{1 + \epsilon \cos \phi} \right\rangle = \frac{a}{(1 - \epsilon^2)^{1/2}} \quad - (17)$$

for all orbitals of H.

$$2) \quad \langle r \rangle = \left\langle \frac{a}{1 + \epsilon \cos \theta} \right\rangle \text{ is different for every orbital of H.}$$

$$3) \quad \langle r \rangle = \left\langle \frac{a}{1 + \epsilon \cos(x\phi)} \right\rangle \text{ is different for all orbitals of H.}$$

$$4) \quad \langle r \rangle = \left\langle \frac{a}{1 + \epsilon \cos(x\theta)} \right\rangle \text{ is different for all orbitals of H.}$$

$$5) \quad \left\langle \frac{1}{r} \right\rangle = \left\langle \frac{1}{a} (1 + \epsilon \cos \phi) \right\rangle \quad - (18)$$

$$= \left\langle \frac{1}{a} (1 + \epsilon \cos \theta) \right\rangle$$

$$= \frac{1}{a}$$

for all orbitals of H.

4)

$$\langle \cos \theta \rangle = \langle \cos \phi \rangle = 0 \quad (19)$$

for all orbitals of H.

Interesting Computations

- 1) $\langle \cos(x\theta) \rangle, \langle \cos(x\phi) \rangle$ for each orbital.
- 2) $\langle \cos(n\theta) \rangle, \langle \cos(n\phi) \rangle$ for each orbital.

The first set corresponds to relativistic correction due to x , and the second set to expectation values for Eckart quantization.
