

280(6): A Photo Theory of Total Internal  
Reflection.

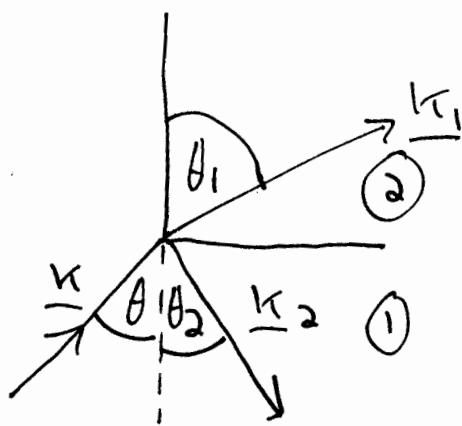
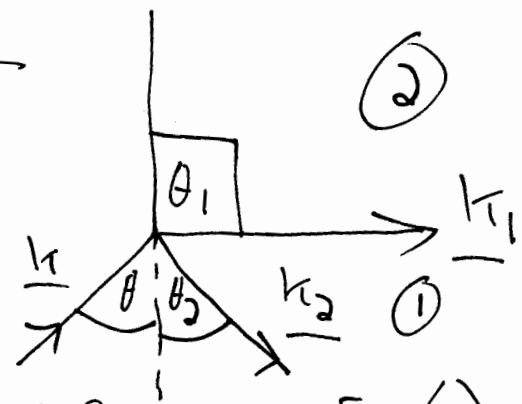


Fig (1)



Total internal reflection is defined in Fig (1),  
if refracted light travels along the boundary  
By conservation of energy and momentum:

$$\langle \omega \rangle = \langle \omega_1 \rangle + \langle \omega_2 \rangle - (1)$$

$$\text{and} \quad \langle \underline{k} \rangle = \langle \underline{k}_1 \rangle + \langle \underline{k}_2 \rangle - (2)$$

Here:

$$\underline{k} = k(i \sin \theta + j \cos \theta) - (3)$$

$$\underline{k}_1 = k_1(i \sin \theta_1 + j \cos \theta_1) - (4)$$

$$\underline{k}_2 = k_2(i \sin \theta_2 - j \cos \theta_2) - (5)$$

It is assumed that medium 1 is glass, in which  
the phase velocity of light is  $v$ , and medium 2,  
air, in which the phase velocity is  $c$ . So

$$k = \frac{\omega}{v}, \quad k_1 = \frac{\omega_1}{c}, \quad k_2 = \frac{\omega_2}{v} - (6)$$

In Boltzmann averaging of the Planck theory:

$$2) \quad \langle \omega \rangle = \left( \frac{\omega}{1-\omega} \right) \omega, \quad \omega = \exp \left( -\frac{\hbar \omega}{kT} \right) \quad (7)$$

$$\langle \omega_1 \rangle = \left( \frac{\omega_1}{1-\omega_1} \right) \omega_1, \quad \omega_1 = \exp \left( -\frac{\hbar \omega_1}{kT} \right) \quad (8)$$

$$\langle \omega_2 \rangle = \left( \frac{\omega_2}{1-\omega_2} \right) \omega_2, \quad \omega_2 = \exp \left( -\frac{\hbar \omega_2}{kT} \right) \quad (9)$$

under the condition:

$$\hbar \omega \ll kT \text{ etc.} \quad (10)$$

therefore:  $\langle \underline{k} \rangle = \left( \frac{\omega}{1-\omega} \right) \frac{\omega}{\sqrt{}} \quad (11)$

$$\langle \underline{k}_1 \rangle = \left( \frac{\omega_1}{1-\omega_1} \right) \frac{\omega_1}{\sqrt{}} \quad (12)$$

$$\langle \underline{k}_2 \rangle = \left( \frac{\omega_2}{1-\omega_2} \right) \frac{\omega_2}{\sqrt{}} \quad (13)$$

From eq. (2):

$$\langle \underline{k}_1 \rangle \cdot \langle \underline{k}_1 \rangle = \langle \underline{k} \rangle \cdot \langle \underline{k} \rangle + \langle \underline{k}_2 \rangle \cdot \langle \underline{k}_2 \rangle - 2 \langle \underline{k} \rangle \cdot \langle \underline{k}_2 \rangle \quad (14)$$

The refractive index of the glass is:

$$n = \frac{c}{\sqrt{}} \quad (15)$$

so:

$$\left( \frac{\omega_1}{1-\omega_1} \right)^2 \omega_1^2 = n^2 \left[ \left( \frac{\omega}{1-\omega} \right)^2 \omega^2 + \left( \frac{\omega_2}{1-\omega_2} \right)^2 \omega_2^2 - 2 \left( \frac{\omega}{1-\omega} \right) \left( \frac{\omega_2}{1-\omega_2} \right) \omega \omega_2 \cos(2\theta) \right] \quad (16)$$

where we have used

$$\theta = \theta_2 - (17)$$

from Snell's Law. So the angle between  $\vec{k}$  and  $\vec{k}_2$  is  $2\theta$ . By Snell's Law for total internal reflection:

$$\begin{aligned} \cos(2\theta) &= \sin^2 \theta - \cos^2 \theta \\ &= 1 - \frac{2}{n^2}. \end{aligned} \quad -(18)$$

Finally, from eq. (1) :

$$\left(\frac{\gamma c}{1-\gamma c}\right)\omega = \left(\frac{\gamma c_1}{1-\gamma c_1}\right)\omega_1 + \left(\frac{\gamma c_2}{1-\gamma c_2}\right)\omega_2 \quad -(19)$$

Eqs. (16) and (19) may be used to eliminate  $\omega_1$  to find  $\omega_2$  in terms of  $\omega$ . From eq. (19) :

$$\left(\frac{\gamma c_1}{1-\gamma c_1}\right)\omega_1 = \left(\frac{\gamma c}{1-\gamma c}\right)\omega - \left(\frac{\gamma c_2}{1-\gamma c_2}\right)\omega_2 \quad -(20)$$

so :

$$\begin{aligned} &\left[ \left(\frac{\gamma c}{1-\gamma c}\right)\omega - \left(\frac{\gamma c_2}{1-\gamma c_2}\right)\omega_2 \right]^2 \\ &n^2 \left[ \left(\frac{\gamma c}{1-\gamma c}\right)^2 \omega^2 + \left(\frac{\gamma c_2}{1-\gamma c_2}\right)^2 \omega_2^2 \right] - (21) \\ &- 2 \left( \frac{\gamma c}{1-\gamma c} \right) \left( \frac{\gamma c_2}{1-\gamma c_2} \right) \omega \omega_2 \left( 1 - \frac{2}{n^2} \right) \end{aligned}$$