

330(4): Systematic Development of Scheme of Quantization
 The classical spin orbit Hamiltonian in the $SU(2)$ basis is:

$$H_{so} = \frac{1}{4m^2 c^2} \underline{\sigma} \cdot \underline{p} \, U \, \underline{\sigma} \cdot \underline{p} \quad - (1)$$

where \underline{p} is the relativistic momentum. Quantization scheme in the Dirac equation is:

$$\underline{p}^\mu = i\hbar \nabla^\mu \quad - (2)$$

and quantization takes place through the relativistic four momentum

$$\underline{p}^\mu = \left(\frac{E}{c}, \underline{p} \right) \quad - (3)$$

This fact almost never appears in the textbooks. Here:

$$E = \gamma mc^2, \quad \underline{p} = \gamma m \underline{v} = \gamma \underline{p}_0 \quad - (4)$$

where \underline{p}_0 is the classical momentum:

$$\underline{p}_0 = m \underline{v} \quad - (5)$$

The Schrodinger equation is not relativistic, so quantization takes place through the classical momentum. The wave functions of the Dirac equation are denoted ψ_r , and those of the Schrodinger equation are denoted ψ . It follows that:

$$-i\hbar \nabla \psi_r = \underline{p} \psi_r \quad - (6)$$

and

$$-i\hbar \nabla \psi = \underline{p}_0 \psi \quad - (7)$$

So in eq. (6) :

$$\underline{p} \phi_r = \underline{p}_0 \phi_r = -i\hbar \underline{\nabla} \phi_r \quad (8)$$

Therefore :

$$H_{so} = \frac{1}{4m^2c^2} \underline{\sigma} \cdot \underline{p}_0 U \underline{\sigma} \cdot \underline{p}_0 \quad (9)$$

$$= \frac{1}{4m^2c^2} \underline{\sigma} \cdot \underline{p} U \underline{\sigma} \cdot \underline{p}_0$$

So :

$$H_{so} \phi_r = -\frac{\hbar}{4m^2c^2} \underline{\sigma} \cdot \underline{\nabla} U \underline{\sigma} \cdot \underline{p}_0 \phi_r \quad (10)$$

$$= -\frac{\hbar}{4m^2c^2} \underline{\sigma} \cdot \underline{\nabla} \left(U \underline{\sigma} \cdot \underline{p}_0 \phi_r \right)$$

as in previous notes for UFT330.

If the relativistic wavefunction ϕ_r is approximated by the Schrodinger wavefunction ϕ then to a first approximation :

$$\phi_r \sim \phi \quad (11)$$

and

$$\underline{p} \phi = \underline{p}_0 \phi \quad (12)$$

So eq. (10) is approximated by :

$$3) \quad H_{so} \psi = - \frac{\hbar \gamma}{4m^2 c^2} \underline{\sigma} \cdot \underline{\nabla} (U \underline{\sigma} \cdot \underline{r} \psi) \quad - (13)$$

The advantage of using the approximation (13) is that the hydrogenic ψ are known analytically.

In these schemes:

$$\gamma = 1 + \frac{H_0 - U}{mc^2} \quad - (14)$$

where

$$U = - \frac{e^2}{4\pi \epsilon_0 r} \quad - (15)$$

Therefore:

$$\gamma = 1 + \frac{H_0}{mc^2} + \frac{e^2}{4\pi \epsilon_0 mc^2 r} \quad - (16)$$

In the H atom:

$$H_0 = \langle H_0 \rangle = - \frac{me^4}{32\pi^2 \epsilon_0^2 \hbar^2 n^2} \quad - (17)$$

where n is the principal quantum number.

So:

$$\gamma = 1 - \frac{e^4}{32\pi^2 \epsilon_0^2 \hbar^2 c^2 n^2} + \frac{e^2}{4\pi \epsilon_0 mc^2 r} \quad - (18)$$

Using the Bohr radius:

4)

$$a_0 = \frac{4\pi\epsilon_0\hbar^2}{me^2} \quad - (19)$$

it follows that

$$H_0 = \langle H_0 \rangle = -\frac{e^2}{8\pi\epsilon_0 mc^2 a_0} \quad - (20)$$

so

$$\gamma = 1 - \frac{e^2}{4\pi\epsilon_0 mc^2} \left(\frac{1}{2a_0} - \frac{1}{r} \right) \quad - (21)$$

A lot of new information appears by careful consideration of γ in various quantization schemes.
Consider the definition of the fine structure constant:

$$d = \frac{e^2}{4\pi\hbar c\epsilon_0}, \quad - (22)$$

then:

$$\gamma = 1 - d \left(\frac{\hbar}{mc} \right) \left(\frac{1}{2a_0} - \frac{1}{r} \right) \quad - (23)$$

i.e.

$$\gamma = 1 - d \lambda_c \left(\frac{1}{2a_0} - \frac{1}{r} \right) \quad - (24)$$

where

$$\lambda_c = \frac{\hbar}{mc} \quad - (25)$$

is the Compton wavelength.

In S.I. units:

5)

$$\alpha = 0.007297351 \quad - (26)$$

$$a_0 = 5.29177 \times 10^{-11} \text{ m} \quad - (27)$$

$$\lambda_c = 3.861591 \times 10^{-13} \text{ m} \quad - (28)$$

Therefore $H_0 = \langle H_0 \rangle = - \frac{\alpha \lambda_c}{2a_0} \quad - (29)$

and $U = - \frac{\alpha \lambda_c}{r} \quad - (30)$

The fundamental approximation used is :

$$\gamma = \left(1 - \frac{p_0^2}{m^2 c^2} \right)^{-1/2} \quad - (31)$$

where $H_0 = \frac{p_0^2}{2m} + U \quad - (32)$

so $p_0^2 = 2m(H_0 - U) \quad - (33)$

The Hamiltonian of the Schrodinger H atom is given by eq. (29) and is negative valued. The reason for this is that H_0 defines binding states between electron and proton, and is defined by :

$$H_0 = \langle H_0 \rangle = - \frac{\hbar^2}{2m} \int \psi^* \nabla^2 \psi d\tau - \frac{e^2}{4\pi\epsilon_0} \int \frac{\psi^* \psi}{r} d\tau \quad - (34)$$

The expectation value of p_0^2 is:

$$p_0^2 = \langle p_0^2 \rangle = -\hbar^2 \int \psi^* \nabla^2 \psi d\tau \quad (35)$$

and is negative valued. This indicates that p_0 is complex valued in the H atom, and this is the result of the Schrodinger equation.

The richness of structure follows from the Leibnitz equation applied to Eq. (10):

$$\begin{aligned} \nabla (\underline{U} \underline{\sigma} \cdot \underline{p} \psi_r) &= \nabla (\underline{\sigma} \cdot \underline{p}) (\underline{U} \psi_r) + \underline{\sigma} \cdot \underline{p} \nabla (\underline{U} \psi_r) \\ &= \nabla (\underline{\sigma} \cdot \gamma \underline{p}_0) (\underline{U} \psi_r) + \underbrace{\gamma \underline{\sigma} \cdot \underline{p}_0 \nabla (\underline{U} \psi_r)}_{\text{Note 330(2)}} \quad (36) \end{aligned}$$

The first term in Eq. (36) can be developed using the Leibnitz theorem:

$$\nabla (\gamma \underline{\sigma} \cdot \underline{p}_0) = (\nabla \gamma) \underline{\sigma} \cdot \underline{p}_0 + \gamma \nabla (\underline{\sigma} \cdot \underline{p}_0) \quad (37)$$

So there is a new term:

$$\hat{H}_{so} \psi = \frac{-i\hbar}{4m^2 c} \left(\underline{\sigma} \cdot \nabla \gamma \underline{\sigma} \cdot \underline{p}_0 \right) \underline{U} \psi_r + \dots \quad (38)$$

1) where:

$$\begin{aligned}\underline{\nabla} \gamma &= - \frac{1}{mc^2} \underline{\nabla} U \\ &= \frac{e^2}{4\pi mc^2 \epsilon_0} \frac{\underline{r}}{r^3}\end{aligned} \quad - (39)$$

giving:

$$\begin{aligned}\hat{H}_{so} \phi_r &= - \frac{i\hbar}{4m^2 c^2} \left(\frac{e^2}{4\pi mc^2 \epsilon_0 r^3} \right) \left(\frac{-e^2}{4\pi \epsilon_0 r} \right) \\ &\times \underline{\sigma} \cdot \underline{r} \underline{\sigma} \cdot \underline{p}_0 \phi_r \quad - (40)\end{aligned}$$

giving:

$$\text{Re} \hat{H}_{so} \phi_r = \frac{-e^4 \hbar}{128 m^3 c^4 \epsilon_0^2} \frac{\underline{\sigma} \cdot \underline{L}}{r^4} \quad - (41)$$
