

398(4): Components of the Spin Connection Vector $\underline{\omega}$ for the Lamb Shift.

In ECE2 theory:

$$\underline{E} = -\underline{\nabla} \phi_0 + \underline{\omega} \phi_0 \quad (1)$$

where \underline{E} is the electric field strength and ϕ_0 the scalar potential. Here $\underline{\omega}$ is the spin connection vector.

$$\underline{\omega} = \omega_x \underline{i} + \omega_y \underline{j} + \omega_z \underline{k} \quad (2)$$

In eq. (1): $\underline{E}(\text{vac}) = \underline{\omega} \phi_0 \quad (3)$

i.e. the vacuum electric field, and:

$$\underline{E}_0 = -\underline{\nabla} \phi_0 \quad (4)$$

i.e. the electric field strength in the hypothetical absence of the vacuum. So:

$$\langle \Delta \underline{E} \rangle = \underline{E} - \underline{E}_0 = \Delta \underline{E} = \underline{\omega} \phi_0 \quad (5)$$

However, the tensor Taylor series gives:

$$\langle \Delta \underline{E} \rangle = \langle \Delta E_x \rangle \underline{i} + \langle \Delta E_y \rangle \underline{j} + \langle \Delta E_z \rangle \underline{k} \quad (6)$$

in which:

$$\langle \Delta E_x \rangle = \langle \Delta E_x \rangle^{(2)} + \langle \Delta E_x \rangle^{(4)} + \langle \Delta E_x \rangle^{(6)} + \dots \quad (7)$$

and similarly for the y and z components. The complete expressions are given by using:

$$f = E_x, E_y, E_z \quad (8)$$

in Eqs. (2) to (4) of Note 398(1). For example at second order in the Taylor series:

$$\langle \Delta E_x \rangle^{(2)} = \frac{1}{6} \langle \underline{S}_i \cdot \underline{S}_j \rangle \nabla^2 E_x \quad - (9)$$

$$E_x = - \frac{e^2 X}{4\pi \epsilon_0 (x^2 + y^2 + z^2)^{3/2}} \quad - (10)$$

Similarly for y and z components.

As in UFT 397:

$$\langle E_x \rangle^{(2)} = \langle E_y \rangle^{(2)} = \langle E_z \rangle^{(2)} = 0 \quad - (11)$$

$$\underline{\omega} \phi_0 = \langle \Delta E_x \rangle \underline{i} + \langle \Delta E_y \rangle \underline{j} + \langle \Delta E_z \rangle \underline{k} \quad - (12)$$

$$-\frac{e^2}{4\pi \epsilon_0 r} \omega_x = \langle \Delta E_x \rangle^{(4)} + \langle \Delta E_x \rangle^{(6)} + \dots \quad - (13)$$

$$-\frac{e^2}{4\pi \epsilon_0 r} \omega_y = \langle \Delta E_y \rangle^{(4)} + \langle \Delta E_y \rangle^{(6)} + \dots \quad - (14)$$

$$-\frac{e^2}{4\pi \epsilon_0 r} \omega_z = \langle \Delta E_z \rangle^{(4)} + \langle \Delta E_z \rangle^{(6)} + \dots \quad - (15)$$

$$\text{where } r = (x^2 + y^2 + z^2)^{1/2} \quad - (16)$$

These are expressions for ω_x , ω_y and ω_z for the Coulomb interaction between the electron and proton in the H atom. The right hand side of eqs. (13) to (15) the fluctuations after isotropic averaging are given by the expressions derived in the previous note:

$$\langle \underline{S}_r \cdot \underline{S}_r \rangle = \frac{2}{\pi} d\lambda^3 \int_{\pi/a_0}^{mc/\hbar} \frac{dk}{k} = \frac{2}{\pi} d\lambda^3 \log_e \frac{1}{\pi d} \quad - (17)$$

$$\langle (\underline{S}_r \cdot \underline{S}_r)^2 \rangle = \frac{4}{\nabla} (d\lambda^3)^2 \int_{\pi/a_0}^{mc/\hbar} \frac{dk}{k^4} \quad - (18)$$

$$\langle (\underline{S}_r \cdot \underline{S}_r)^3 \rangle = \frac{8\pi}{\nabla^2} (d\lambda^3)^3 \int_{\pi/a_0}^{mc/\hbar} \frac{dk}{k^7} \quad - (19)$$

The Lamb shift is:

$$\begin{aligned} \langle \Delta \phi \rangle^{(2)} &= \frac{1}{6} \langle \underline{S}_r \cdot \underline{S}_r \rangle \frac{e^2}{\epsilon_0} |\psi(0)|^2 \quad - (20) \\ &= \frac{\hbar c}{6\pi a_0} \log_e \frac{1}{\pi d} \end{aligned}$$

and must be calculated with:

$$\nabla^2 \left(\frac{1}{r} \right) = -4\pi \delta_D \quad - (21)$$

where δ_D is the Dirac delta function.

Note carefully that a classical level of v_r (13) to (15):

$$\nabla^2 \left(\frac{1}{r} \right) = 0 \quad - (22)$$

and

$$\nabla^2 \left(\frac{1}{r^2} \right) = 0 \quad - (23)$$

so v_r (11) follows.