

5(1): Precession of a Satellite such as Gravity Probe B

1. Thomas Precession

Consider the  $E(2)$  covariant metric in plane polar coordinates  $(r, \phi)$ :

$$ds^2 = c^2 d\tau^2 = c^2 dt^2 - dr^2 - r^2 d\phi^2 \quad (1)$$

Thomas precession is due to the frame rotation:

$$\phi' = \phi + \omega t \quad (2)$$

in UFT110, a heavily studied paper that appears on the first page of Google with keywords "Thomas precession derivation". It follows from eq. (2) that:

$$d\phi' = d\phi + \omega dt \quad (3)$$

and

$$d\phi'^2 = d\phi^2 + 2\omega d\phi dt + \omega^2 dt^2 \quad (4)$$

so

$$ds'^2 = (c^2 - r^2 \omega^2) dt^2 - 2\omega r^2 d\phi dt - dr^2 - r^2 d\phi^2 \quad (5)$$

The orbit of Gravity Probe B is nearly circular, so:

$$v = \omega r \quad (6)$$

where  $v$  is the orbital velocity, so:

$$ds'^2 = \left(1 - \frac{v^2}{c^2}\right) \left(c^2 dt^2 - 2r^2 \Omega d\phi dt\right) - dr^2 - r^2 d\phi^2 \quad (7)$$

where

$$\Omega = \omega \left(1 - \frac{v^2}{c^2}\right)^{-1} \quad (8)$$

is the relativistic angular velocity.

The infinitesimal of time changes to:

$$dt' = \left(1 - \frac{v^2}{c^2}\right)^{-1/2} dt - (9)$$

Therefore is a rotation of  $2\pi$ , the Thomas precession

$$\Delta\phi = \Omega dt' - \omega t$$

$$= 2\pi \left( \left(1 - \frac{v^2}{c^2}\right)^{-1/2} - 1 \right) - (10)$$

For the approximately circular orbit of GPR:

$$v^2 = \frac{MG}{r} - (11)$$

where  $M$  is the mass of the Earth, and  $r$  is the distance of Gravity Probe B to the centre of the earth:

$$M = 5.98 \times 10^{24} \text{ kg} - (12)$$

$$G = 6.67 \times 10^{-11} \text{ m}^3 \text{ kg}^{-1} \text{ s}^{-2} - (13)$$

$$r = 7.0274 \times 10^6 \text{ m} - (14)$$

So the Thomas precession of GPR is:

$$\Delta\phi = 2\pi \left( \left(1 - \frac{MG}{rc^2}\right)^{-1/2} - 1 \right) - (15)$$

$$\sim \pi \frac{MG}{rc^2}$$

The orbit of Gravity Probe B takes 90 minutes to complete.

Lense-Thirring Precession at the Equator.

At the equator, the magnitude of the gravitomagnetic field is:

$$\Omega = \frac{2}{5} \frac{MG R^2 \omega_E}{c^2 r^3} - (16)$$

where  $R$  is the radius of the Earth,  $r$  is given by eq. (14), and  $\omega_E$  is the angular velocity of the Earth. At - UFT 110 :

$$\Omega = 1.52 \times 10^{-14} \text{ rad s}^{-1} \quad - (17)$$

The Lense Thirring precession per orbit is

$$\Delta \phi_{LT} = \frac{1}{2} \Omega t \quad - (18)$$

also  $t = 90 \text{ minutes} = 5,400 \text{ seconds} \quad - (19)$

so  $\Delta \phi_{LT} = 4.10 \times 10^{-11} \text{ radians per orbit} \quad - (20)$

The Lense Thirring precession is roughly two orders of magnitude smaller than the Thomas precession.

### 3) Einsteinian orbital precession

This is  $\Delta \phi_E = \frac{3\pi G}{ac^2(1-e^2)} \quad - (21)$

from the method of Meria and Thornton, and

$$\Delta \phi_E = \frac{3\pi G}{ac^2(1-e)} \quad - (22)$$

from the apsidal method for an orbit of  $2\pi$  radians. The eccentricity of Gravity Probe B is :

$$e = 0.0014 \quad - (23)$$

$$a \approx r \quad - (24)$$

so  $\Delta \phi_E = 1.894 \times 10^{-9} \text{ radians per orbit.}$

4) Note that for a nearly circular orbit the Thomas precession is

$$\Delta\phi_T = \frac{\pi \hbar G}{rc^2} - (17)$$

and the Einsteinian precession is

$$\Delta\phi_E = \frac{3\hbar G}{rc^2} - (18)$$

and they are almost the same, because:

$$\pi = 3.1415927 - (19)$$

It is well known that the Einstein theory is totally incorrect, as is the classic UFT88, so. planetary precession is Thomas precession, a correct theory. Thomas precession is correctly E(2) covariant.

The next note will evaluate the spin corrections for these precessions, and consider geodesic precession using the method of UFT345.

Frankly, Mike B can be severely criticized on the grounds that it claims to see a small precession, the Lense Thirring precession, but this is two orders of magnitude smaller than the Thomas precession. The standard model never considers the Thomas precession of planets.