

# Paper 61 : Summary of the Coulomb Law in ECE Theory

The law is given from the first Cartan structure equation:

$$\tau^a = d \wedge q^a + \omega^a_b \wedge q^b \quad - (1)$$

and the first Bianchi identity:

$$d \wedge \tau^a + \omega^a_b \wedge \tau^b = R^a_b \wedge q^b, \quad - (2)$$

w/ the ECE Ansatz:

$$\left. \begin{aligned} A^a &= A^{(0)} q^a \\ F^a &= A^{(0)} \tau^a \end{aligned} \right\} - (3)$$

Eqs (1) to (3) lead to:

$$\underline{\nabla} \cdot \underline{E}^a = c \mu_0 \tilde{J}^{0a} \quad - (4)$$

$$\underline{E}^a = - \frac{\partial \underline{A}^a}{\partial t} - \underline{\nabla} \phi^a - c \omega^{0a}_b \underline{A}^b + \underline{\omega}^a_b \phi^b \quad - (5)$$

Eq. (4) may be written for each index  $a$  as:

$$\underline{\nabla} \cdot \underline{E} = \rho / \epsilon_0 \quad - (6)$$

If we restrict attention to the scalar potential then for each index  $a$ , eq. (5) is:

$$\underline{E} = - \underline{\nabla} \phi + \underline{\omega}_b \phi^b \quad - (7)$$

In eq (7)  $\phi^b$  is interpreted as a scalar quantity indexed by a label  $b$ . For a particular index therefore, eq. (7) is:

$$\underline{E} = - \underline{\nabla} \phi + \underline{\omega} \phi \quad - (8)$$

2) Eqs (6) and (8) give the equation:

$$\underline{\nabla} \cdot (-\underline{\nabla} \phi + \underline{\omega} \phi) = \rho / \epsilon_0 \quad - (9)$$

i.e

$$\boxed{\nabla^2 \phi - \underline{\nabla} \cdot (\phi \underline{\omega}) = -\rho / \epsilon_0} \quad - (10)$$

The spin connection vector is:

$$\left. \begin{aligned} \underline{\omega} &= \omega_x \underline{i} + \omega_y \underline{j} + \omega_z \underline{k} \\ &= \omega_r \underline{e}_r + \omega_\phi \underline{e}_\phi + \omega_\theta \underline{e}_\theta \end{aligned} \right\} - (11)$$

where  $\omega_r$  is the radial component of  $\underline{\omega}$ . If the latter is purely radial, then:

$$\underline{\omega} = \omega_r \underline{e}_r \quad - (12)$$

and:

$$\underline{\omega} \cdot \underline{\nabla} \phi = \omega_r \frac{\partial \phi}{\partial r} \quad - (13)$$

$$\phi \underline{\nabla} \cdot \underline{\omega} = \frac{\phi}{r^2} \frac{\partial}{\partial r} (r^2 \omega_r) \quad - (14)$$

$$\nabla^2 \phi = \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial \phi}{\partial r} \right) = \frac{\partial^2 \phi}{\partial r^2} + \frac{2}{r} \frac{\partial \phi}{\partial r} \quad - (15)$$

It is reasonable to assume that:

$$\omega_r = \frac{A}{r} \quad - (16)$$

where  $A$  is dimensionless. Then eqs. (10) to (16)

give the result:

$$\frac{d^2 \phi}{dr^2} + \frac{(2-A)}{r} \frac{d\phi}{dr} - \frac{A\phi}{r^2} = -\frac{\rho}{\epsilon_0} \quad (17)$$

This equation has a similar structure to the well known one-dimensional Schrödinger equation:

$$\frac{d^2 P}{dr^2} + \frac{2\mu}{\hbar^2} \frac{e^2}{4\pi\epsilon_0} \frac{P}{r} - \frac{2\mu}{\hbar^2} \frac{l(l+1)\hbar^2}{2\mu r^2} P = -\frac{2\mu E}{\hbar^2} P \quad (18)$$

for which is the effective potential:

$$V_{\text{eff}} = -\frac{e^2}{4\pi\epsilon_0 r} + \frac{l(l+1)\hbar^2}{2\mu r^2} \quad (19)$$

However, eq. (17) contains second and first r partial derivatives of  $\phi$ , and a term in  $\phi$ .

The comparison of structure of eqs. (17) and (18) shows that the term in  $-\frac{A\phi}{r^2}$  is repulsive in nature.

If  $A = 2 \quad (20)$

eq. (17) becomes:

$$\frac{d^2 \phi}{dr^2} - \frac{2\phi}{r^2} = -\frac{\rho}{\epsilon_0} \quad (21)$$

4) The ~~Laplace~~ <sup>Poisson</sup> equation is:

$$\nabla^2 \phi = -\frac{\rho}{\epsilon_0} \quad - (22)$$

and this describes the attractive part of the Coulomb law. If  $\rho$  is very small, eq. (21) becomes:

$$\frac{d^2 \phi}{dr^2} \sim \frac{2\phi}{r^2} \quad - (23)$$

so:

$$\phi = \alpha r^2 + \frac{\beta}{r} \quad - (24)$$

The Coulomb potential <sup>energy</sup> is:

$$V_c = -e \phi_c = -\frac{e^2}{4\pi \epsilon_0 r} \quad - (25)$$

The effect of the spin concentration is to add a repulsion of approximate form (24) to the attractive Coulomb potential (25). It is seen from eq. (21) that:

$$\frac{d^2 \phi}{dr^2} = \frac{2\phi}{r^2} - \frac{\rho}{\epsilon_0} \quad - (26)$$

so as  $r \rightarrow 0$ ,  $\phi \rightarrow \infty$ , and the repulsion overleaps the attraction.

For one electron a positive energy

5) is inputted of  $P_0$  form:

$$\bar{V} = e \left( \alpha r^2 + \frac{\beta}{r} \right) \quad - (27)$$

and if  $P_0$  exceeds 13.6 eV of H atom is ionized.

Eq (27) is a general law and shows that it is possible to input energy from fracture into a material to release free electron. The mechanism to do this must be found experimentally.

At macroscopic distances in laboratory experiments the term  $\alpha/r^2$  is very small compared with  $-\beta/r$ , and the Coulomb law applies.

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