

4(i) : ECE Dynamics of a Massive Spin-1 Particle

In the standard model these dynamics are not gauge invariant, so violate the gauge principle. In ECE theory the problem can be developed straight forwardly. This is illustrated with the massive photon. In the

limit:

$$kT = \left(\frac{mc}{\hbar} \right)^2 \quad \text{--- (1)}$$

The wave equation of the massive photon is:

$$\left(\square + \left(\frac{mc}{\hbar} \right)^2 \right) A_\mu^a = 0 \quad \text{--- (2)}$$

which is the Proca equation for each a , i.e. for each polarization:

$$a = (0), (1), (2), (3) \quad \text{--- (3)}$$

The e/m field is:

$$F_{\mu\nu}^a = \partial_\mu A_\nu^a - \partial_\nu A_\mu^a + \omega_{\mu\nu}^a A_\nu^b - \omega_{\nu\mu}^a A_\mu^b \quad \text{--- (4)}$$

Thus:

$$\partial^\mu F_{\mu\nu}^a = \square A_\nu^a + \partial^\mu (\omega_{\mu\nu}^a A_\nu^b - \omega_{\nu\mu}^a A_\mu^b) \quad \text{--- (5)}$$

Eq. (2) is the equation of the free photon of mass m . (Comparing eqs (2) and (5):

$$\left(\frac{mc}{\hbar}\right)^2 A_{\sim}^a = \mathcal{J}^{\mu} \left(\omega_{\mu b}^a A_{\sim}^b - \omega_{\sim b}^a A_{\mu}^b \right) \quad - (6)$$

This is a free space equation, because we are considering a free photon, i.e. are free of ~~the~~ gravitational field of matter, such as an electron.

Now use:

$$A_{\sim}^a A_{\sim}^a = 4 A^{(0)2} \quad - (7)$$

It is found:

$$m^2 = \frac{\hbar^2}{4c^2 A^{(0)2}} A_{\sim}^a \mathcal{J}^{\mu} \left(\omega_{\mu b}^a A_{\sim}^b - \omega_{\sim b}^a A_{\mu}^b \right) \quad - (8)$$

Also from eq. (6):

$$\left(\frac{mc}{\hbar}\right)^2 \mathcal{J}^{\sim} A_{\sim}^a = 0 \quad - (9)$$

so

$$\mathcal{J}^{\sim} A_{\sim}^a = 0 \quad - (10)$$

From eq. (8) the photon mass is:

$$m \sim 10^{-42} \text{ kgms} \quad - (11)$$

3) From eq. (10) we see that the Lorentz
condition always holds. From eq. (8) the photon
mass is due to the spin connection.

In the presence of matter:

$$\left(\frac{mc}{\hbar}\right)^2 \rightarrow k_T \quad - (12)$$

and
$$g^{\mu\nu} F_{\mu\nu}^a \neq 0 \quad - (13)$$

This introduces the inhomogeneous current, so let

$$g^{\mu\nu} F_{\mu\nu}^a = \mu_0 J_{\nu}^a \quad - (14)$$

So J_{ν}^a is known to the difference between
 k_T and $\left(\frac{mc}{\hbar}\right)^2$.

Standard Model

The Proca equation is:

$$F_{\mu\nu} = \partial_{\mu} A_{\nu} - \partial_{\nu} A_{\mu} \quad - (15)$$

$$d_{\mu} F^{\mu\nu} + \left(\frac{mc}{\hbar}\right)^2 A^{\nu} = 0 \quad - (16)$$

So:

$$\left(\frac{mc}{\hbar}\right)^2 \partial_{\nu} A^{\nu} = 0 \quad - (17)$$

4) and $\left(\mathbb{1} + \left(\frac{mc}{\hbar} \right)^2 \right) A_\mu = 0 \quad - (18)$

The Lagrangian for the Proca equation is:

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \frac{1}{2} m^2 A_\mu A^\mu \quad - (19)$$

(reduced units).

This poses a fatal weakness for the standard model because neither eq. (19) nor eq. (17) is gauge invariant. A massive spin one particle cannot be understood with a gauge principle.

In ECE this is not a problem because the gauge principle is not used. In ECE the origin of photon mass is traced to the spin connection.
