REFUTATION OF THE CLAIMS OF G. BRUHN he relevant e wail posting by G. Drahn is appelled. Bruhn posted his without my howwledge, and appeartly it is an unederseed document. Drahn is dero to none of my recent work which is now universally acclaimed and accepted. Here I comment universally acclaimed and accepted. Here I comment or Draha's posting rection by section. this is nevely a quote by bruke of some of my ealiest work on B(31) almost a decade ago. Section 2 The first elementary error is brukers claims occurs lis kis equalitions (2,3) and (2,4). It is easy to show that bruke has made an elementary blueder be came when we and his ego (2,3) and (2,4) we obtain: B(1)(lest) + B(1)(right) = B(0) e(1)(eit+e-it) $=\frac{B^{(0)}(\underline{i}-\underline{i})(e^{-i\phi})-(\underline{i})}{\sqrt{2}}$ $=\frac{B^{(0)}(i-ij)}{(os\phi)}$ House, it is well known that the sun of left and right circular plaination is linear plaination. Druhe's equ (3) is still crowner plaination

The correct way to represent left and right creater pherization is a follows:

1 (1) = 1 (1) (1) (2) (3) $\underline{B}(i) = \frac{1}{\sqrt{3}}(\underline{B} \times \underline{i} + i\underline{B} + \underline{i})e^{-i\phi} - (4)$ Add (3) and (4) to obtain: $B_{L}^{(1)} + B_{R}^{(1)} = \frac{2}{\sqrt{5}} D_{X} \stackrel{!}{=} e^{-} (5)$ Which is linear placeration as required. Just I fewe nentioned the chirolity

apurtin is my writings many times (e.s. Adv. Chen.

Phys. vl 85). This alway Oat Bruha is almost

constelly grownt of my work. The vest of Bruha's

constelly grownt of my work. The vest of Bruha's

claim is sequentially erroreus, and furthermore, contains

other eleventary bluders. In Ris Section, Bruhn gnotes my well-from I Cyclic Theorem. from eq. (3) un see Part to left (P wave is $B_{L}^{(1)} = \frac{1}{\sqrt{3}} (B_{X} \cos \phi i + B_{Y} \sin \phi j) - (6)$ Section 4

3) For egr. (4) we see that the right (P warn is $\underline{B}_{R} = \frac{1}{5} \left(B_{\times} \cos \phi - B_{Y} \sin \phi \right) - (7)$ Eqn. (6) is Bouhr's eqn. (4.2) and eqn. (7) is Inhis eq. (4.13). This shows that my eans
(3) and (4) give bruke's own definitions of left
and right harder CP. it his eggs (4.1) and (4.2). Brohn's egrs (4.1) and (4.2) Derefre contradict Boshi, egrs (2.3) and (2.4), became when we add bulis eggs (4.1) and (4'2) we obtain linear Merization, Sut ules us add egno (2.3) and (3.4) we obtail a course present in.

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These errors already nain error is to confine underseed claim, but his main error is to confine underseed claim, but his main error is to confine underseed is nearly in left and right of the present of the lis reason that I (3) is egged and opposite for left and right circular plantation, and for this reason plantation, and for this reason plantation. The lend is a like the large plantation.

There is the large plantation of the lend is all your large. Lie office excuer previsation. $\frac{B}{R} = -\frac{i}{B^{(0)}} \frac{B}{R} \times \frac{B}{R} = -\frac{i}{B^{(0)}} \frac{i}{i} \frac{i}{i} \frac{i}{i} \frac{k}{i} - (8)$ $= i \frac{3}{B^{(0)}} \frac{k}{k} = -\frac{B^{(0)}}{B^{(0)}} \frac{k}{k}$ $= -\frac{i}{\mathcal{B}^{(0)}} \underbrace{\mathcal{B}^{(1)}_{1} \times \mathcal{B}^{(2)}_{1}}_{(0)} = -\frac{i}{2} \underbrace{\mathcal{B}^{(0)}_{0}}_{(0)} \begin{vmatrix} i & j & k \\ 1 & -i & 0 \\ 1 & i & 0 \end{vmatrix} - (9)$

 $\frac{1}{B} = -i^{2} C^{(0)} = B^{(0)} = B^{(0)}$ Thus: $\left| \frac{B}{R} \right| = -B \left(\frac{3}{3} \right)$ - (10) The glaring error in bruhers claims occurs in his equ (4.5), where he asserts that for lular plaisation: $\underline{\mathbb{B}^{(3)}} = ? \pm \sqrt{2} \mathbb{B}^{(0)} \mathbb{k}$ - (II) The regra for the emoneum exp. (11) appears to be confusion about that is ment by B (3). Relatter is usell-defined it my work as existing for one ploton. One ploton is effect left handely, is which (ax v. ostail eg (9), a right hearted, is vhick (on ve ostar egn (8). If we superimore one left hurled flora with one right femder plota we obtain $\underline{B}^{(3)} = \underline{B}^{(3)} + \underline{B}^{(3)} = \underline{O}$. Re \underline{B} (yelic Heven always applies to one serve of presidentia. Renfre Le clair by bruke is false and workstands, i.e. it is made it ignorance of to li Verature.

Muc.

Refutation of Myron W. Evans' B⁽³⁾ field hypothesis

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Summary. In 1992 Myron W. Evans published a paper [1] where he proposed the hypothesis that each circularly polarized plane electromagnetic wave – in addition to its Maxwellian transversal components – should have a longitudinal component of magnitude $B^{(3)} = B_0/2^{1/2}$ compared with the real magnetic flux amplitude B_0 of the circularly polarized plane wave. Two years later, in a paper [2] he added so-called "cyclical relations" that should hold between the components of the flux **B** relative to a certain complex basis $e^{(1)}$, $e^{(2)}$, $e^{(3)}$ for general plane waves. By application to the superposition of two circularly polarized plane waves to a linearly polarized wave we show here that Evans' "cyclical relations" cannot hold generally. The assumption of a longitudinal $B^{(3)}$ field leads to a contradiction. This affects especially the paper [4], where a kind of PMM, the MEG, is justified by means of the $B^{(3)}$ field.

1. Evans' circular basis (taken from [3], p. 7-14, with slight corrections)

Let (x,y,z) denote Cartesian coordinates with unit vectors \mathbf{i} , \mathbf{j} , \mathbf{k} parallel to the corresponding axes. Evans assumes the z-direction to be the direction of propagation of a plane electromagnetic wave. The letter i denotes the imaginary unit as usual.

Evans replaces the Cartesian unit vectors i, j, k with another system of unit vectors called circular basis

(1.1)
$$\mathbf{e}^{(1)} = (\mathbf{i} - \mathbf{i} \mathbf{j})/2^{\frac{1}{2}} \qquad \mathbf{e}^{(2)} = (\mathbf{i} + \mathbf{i} \mathbf{j})/2^{\frac{1}{2}} \qquad \mathbf{e}^{(3)} = \mathbf{k}$$

This means that a certain unitary coordinate transform is executed.

We suppose the coordinates a_x , a_y , a_z of all vectors $\mathbf{a} = a_x \mathbf{i} + a_y \mathbf{j} + a_z \mathbf{k}$ to be real. Then from

$$\mathbf{a_x} \mathbf{i} + \mathbf{a_v} \mathbf{j} + \mathbf{a_z} \mathbf{k} = \mathbf{a} = \mathbf{a}^{(1)} \mathbf{e}^{(1)} + \mathbf{a}^{(2)} \mathbf{e}^{(2)} + \mathbf{a}^{(3)} \mathbf{e}^{(3)}$$

we obtain the transformation rule for coordinates

(1.2)
$$a^{(1)} = 2^{-\frac{1}{2}} (a_x + i a_y), \qquad a^{(2)} = 2^{-\frac{1}{2}} (a_x - i a_y), \qquad a^{(3)} = a_z.$$

Evidently the coordinates fulfil the equation

(1.3)
$$|\mathbf{a}^{(1)}|^2 + |\mathbf{a}^{(2)}|^2 + |\mathbf{a}^{(3)}|^2 = \mathbf{a_x}^2 + \mathbf{a_x}^2 + \mathbf{a_x}^2 = |\mathbf{a}|^2.$$

Additionally the vector components of a relative to the circular basis are defined by

(1.4)
$$\mathbf{a}^{(1)} = \mathbf{a}^{(1)} \mathbf{e}^{(1)}, \qquad \mathbf{a}^{(2)} = \mathbf{a}^{(2)} \mathbf{e}^{(2)}, \qquad \mathbf{a}^{(3)} = \mathbf{a}^{(3)} \mathbf{e}^{(3)}.$$

Let ...* denote the conjugate complex of the term where * is attached. Then evidently we have the symmetry properties

(1.5)
$$\mathbf{e}^{(1)*} = \mathbf{e}^{(2)}, \quad \mathbf{e}^{(2)*} = \mathbf{e}^{(1)}, \quad \mathbf{e}^{(3)*} = \mathbf{e}^{(3)}$$

(1.6)
$$\mathbf{a}^{(1)*} = \mathbf{a}^{(2)}, \quad \mathbf{a}^{(2)*} = \mathbf{a}^{(1)}, \quad \mathbf{a}^{(3)*} = \mathbf{a}^{(3)}$$

and

(1.7)
$$a_1 = a_2^*, a_2 = a_1^*, a_3 = a_3^*.$$

By direct calculation one can obtain the cyclic cross product rules

(1.8)
$$\mathbf{e}^{(1)} \times \mathbf{e}^{(2)} = \mathbf{i} \ \mathbf{e}^{(3)} *, \qquad \mathbf{e}^{(2)} \times \mathbf{e}^{(3)} = \mathbf{i} \ \mathbf{e}^{(1)} *, \qquad \mathbf{e}^{(3)} \times \mathbf{e}^{(1)} = \mathbf{i} \ \mathbf{e}^{(2)} *.$$

2. Evans B⁽³⁾ hypothesis from 1992

In 1992 Myron W. Evans published a paper [1] where he proposed the hypothesis that a monochrome circularly polarized plane electromagnetic wave should have - in addition to its Maxwellian transversal components - a longitudinal component the size of which will be specified at the end of this section.

We assume the speed c of light to be 1. Then a monochrome Maxwellian circularly polarized plane

wave propagating in z-direction with amplitude $B_0 > 0$ is given by the equations

(2.1)
$$B_x = B_0 \cos \omega(t-z), \qquad B_y = \pm B_0 \sin \omega(t-z), \quad B_z = 0$$

Here the sign \pm in B_y determines the chirality of the polarization: The + sign is valid for left circular polarization and the - sign for right circular polarization. Introducing the abbreviation

(2.2)
$$B^{(0)} = 2^{-1/2} B_0$$

the components of the Maxwellian B relative to Evans' circular basis can be written as

(2.3)
$$\mathbf{B}^{(1)} = \mathbf{e}^{(1)} \mathbf{B}^{(0)} \mathbf{e}^{i\omega(t-z)}, \qquad \mathbf{B}^{(2)} = \mathbf{e}^{(2)} \mathbf{B}^{(0)} \mathbf{e}^{-i\omega(t-z)}$$

in case of left circular polarization and

(2.4)
$$\mathbf{B}^{(1)} = \mathbf{e}^{(1)} \mathbf{B}^{(0)} \mathbf{e}^{-i\omega(\mathbf{t}-\mathbf{z})}, \qquad \mathbf{B}^{(2)} = \mathbf{e}^{(2)} \mathbf{B}^{(0)} \mathbf{e}^{i\omega(\mathbf{t}-\mathbf{z})}$$

in case of right circular polarization.

The hypothesis of Evans' paper [1] is that a monochrome circularly polarized plane electromagnetic wave should have - in addition to its Maxwellian transversal components - a longitudinal component of magnitude $B^{(3)} = B^{(0)}$, i.e.

(2.5)
$$\mathbf{B}^{(3)} = \mathbf{e}^{(3)} \mathbf{B}^{(0)}.$$

He does not mention whether there should be a sign dependency on the chirality of the circularly polarized wave.

In summary may be said that the equations (2.3-5) describe the Evans version of a circularly polarized wave, while for the Maxwellian circularly polarized wave equation (2.5) has to be replaced with $\mathbf{B}^{(3)} = \mathbf{0}$.

3. Evans' Cyclic Relations

In 1994 Evans supplemented his former hypothesis from 1992 by another paper [2]. Here he starts with the statement that the magnetic flux vector **B** of each circularly polarized plane wave that he had equipped in [1] with the additional longitudinal component (2.5) fulfils the "cyclic relations"

(3.1)
$$\mathbf{B}^{(1)} \times \mathbf{B}^{(2)} = \mathbf{i} \ \mathbf{B}^{(0)} \mathbf{B}^{(3)*},$$

(3.2)
$$\mathbf{B}^{(2)} \times \mathbf{B}^{(3)} = \mathbf{i} \ \mathbf{B}^{(0)} \mathbf{B}^{(1)*},$$

(3.3)
$$\mathbf{B}^{(3)} \times \mathbf{B}^{(1)} = i \mathbf{B}^{(0)} \mathbf{B}^{(2)*},$$

which can be confirmed easily by means of the equations (2.2-5).

Evans' new hypothesis of 1994 generalizes the equations (3.1-3) to general waves in vacuo [2, p. 69]:

"We assert therefore that in classical electrodynamics there are three components $\mathbf{B}^{(1)}$, $\mathbf{B}^{(2)}$ and $\mathbf{B}^{(3)}$ of a travelling plane wave in vacuo. These are interrelated in the circular basis by equations (3.1-3). The third component, the ghost field

$$\mathbf{B}^{(3)} = \mathbf{B}^{(1)} \times \mathbf{B}^{(2)} / (i \mathbf{B}^{(0)}) = \mathbf{B}^{(0)} \mathbf{k}$$

is real and independent of phase."

Hence Evans' cyclic equations should be valid for the superposition of circularly polarized plane waves too. This is it what we will check now.

4. Superposition of circularly polarized waves

The superposition of a right circularly polarized wave with its left circularly polarized counterpart yields linearly polarized plane waves. If we superpose the right circularly polarized wave

(4.1)
$$\mathbf{B}_{\mathbf{r}} = \mathbf{B}_{0} \left[\mathbf{i} \cos \omega (\mathbf{t} - \mathbf{z}) - \mathbf{j} \sin \omega (\mathbf{t} - \mathbf{z}) \right]$$

and the left circularly polarized wave

(4.2)
$$\mathbf{B}_1 = \mathbf{B}_0 \left[\mathbf{i} \cos \omega (t-z) + \mathbf{j} \sin \omega (t-z) \right],$$

we obtain the linearly polarized wave

$$\mathbf{A}(4.3) \qquad \mathbf{B} = 2\mathbf{B}_0 \mathbf{i} \cos \omega (\mathbf{t} - \mathbf{z}),$$

i.e.
$$B_x = 2I$$

$$B_x = 2B_0 \cos \omega(t-z), \quad B_y = B_z = 0.$$

Due to Evans both circularly polarized waves should be accompanied by ghost fields $\mathbf{B}_r^{(3)}$ and $\mathbf{B}_l^{(3)}$ which give the resulting sum field

(4.4)
$$\mathbf{B}^{(3)} = \mathbf{B_r}^{(3)} + \mathbf{B_l}^{(3)}$$

But due to the indeterminacy of the sign of the additional Evans field for circularly polarized waves we have to discuss all combinations of signs: the cases of constructive and destructive superposition of the corresponding Evans fields. The resulting Evans field for linearly polarized plane waves could be

(4.5)
$$\mathbf{B}^{(3)} = \mathbf{0}$$
 or $\mathbf{B}^{(3)} = \pm 2^{1/2} \mathbf{B}_0 \mathbf{k}$.

We have to check whether there is a combination that fulfils the cyclic equations (3.1-3):

Due to the rules (1.2) the linearly polarized wave (4.3) yields

(4.6)
$$B^{(1)} = 2^{1/2} (B_x + i B_y) = 2^{1/2} B_0 \cos \omega (t-z), \qquad B^{(2)} = B^{(1)*} = 2^{1/2} B_0 \cos \omega (t-z).$$

Hence we get:

Case
$$B^{(3)} = 0$$

Then the equation (3.1) leads to a contradiction, since we obtain

(4.7)
$$\mathbf{B}^{(1)} \times \mathbf{B}^{(2)} = \mathbf{B}^{(1)} \mathbf{B}^{(2)} \mathbf{e}^{(1)} \times \mathbf{e}^{(2)} = 2 \mathbf{B}_0^2 \cos^2 \omega (t-z) \mathbf{i} \mathbf{k} \neq \mathbf{0} = \mathbf{i} \mathbf{B}^{(0)} \mathbf{B}^{(3)}.$$

Cases
$$B^{(3)} = \pm 2^{1/2} B_0$$

Each of the equations (3.2) and (3.3) yields $B^{(0)} = B^{(3)}$. Therefore the right side of (3.1) gives

(4.8)
$$i B^{(0)} B^{(3)*} = i B^{(3)2} k = 2 i B_0^2 k$$

But for the left side of (3.1) we obtain

(4.9)
$$\mathbf{B}^{(1)} \times \mathbf{B}^{(2)} = \mathbf{B}^{(1)} \mathbf{B}^{(2)} \mathbf{e}^{(1)} \times \mathbf{e}^{(2)} = 2 \mathbf{B}_0^{(2)} \cos^2 \omega(t-z) \mathbf{i} \mathbf{k}$$

Hence we have a contradiction again.

Thus Evans' cyclic equations do not hold for the superposition of two circularly polarized plane waves to a linearly polarized plane wave.

Hence the validity of Evans' cyclic equations, of the base of Evans' O(3) theory, is refuted.

References

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