# DEVELOPMENT OF m THEORY PART ONE: RELATIVISTIC KINETIC ENERGY, EINSTEIN ENERGY EQUATION, AND POTENTIAL ENERGY DUE TO m (r).

by

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#### ABSTRACT

The systematic development is initiated of classical dynamics in the most general spherically symmetric spacetime (m space). The Einstein energy equation, the rest energy, the relativistic kinetic energy and the potential energy due to m (r) (the energy due to spacetime, vacuum or aether) are derived systematically. Rigorous self consistency of concepts and methods is demonstrated.

Keywords: Development of m theory, relativistic kinetic energy, Einstein energy equation, rest energy, potential energy from m (r).

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### 1. INTRODUCTION

In the immediately preceding paper (UFT417) of this series {1 - 41}, the natural philosophy (m theory) was initiated of the most general spherically symmetric spacetime (m space). It was shown that in that type of naturally philosophy or physics. superluminal signalling, infinite energy from m space, and counter gravitation are rigorously possible. In Section 2 of this paper the m theory is developed systematically in classical dynamics to derive the relativistic kinetic energy the Einstein energy equation, the rest energy, and potential energy from m (r). The latter is defined by the infinitesimal line element of m space. The rigorous self consistency of concepts and methods is demonstrated in several complementary ways.

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This paper is a brief synopsis of extensive calculations found in the notes accompanying UFT418 on www.aias.us. Note 418(1) defines the classical limit of the hamiltonian and lagrangian of m theory, which is part of the ECE and ECE2 generally covariant unified field theory. Note 418(2) is a demonstration of the rigorous self consistency of the Hamiltonian and Lagrangian methods used in m theory. Note 418(3) the flat spacetime limit, the work integral and begins the derivation of the relativistic kinetic energy in m space. Note 418(4) derives the relativistic kinetic energy of m space from the work integral over the relativistic momentum of m space. Note 418(5) is a convenient summary of m space classical dynamics, a demonstration of the rigorous self consistency of the Lagrangian method used in m space, a derivation of the ECE2 spin connection from m theory and initial development of a concept unique to m theory: the potential energy of m space. This is the rigorous definition of "energy from spacetime", "vacuum energy", or "aether energy". Note 418(6) derives the Einstein energy equation of m space from its relativistic linear momentum. This method also derives the rest energy of m space. Notes 418(7) and 418(8) give a rigorously self outnisiste derivation of the important potential energy due to m space ("the energy of the vacuum").

In Section 3, the calculations and concepts of Section 2 are analysed numerically and graphically, so that the meaning of the complicated mathematics becomes clear.

#### 2. FUNDAMENTAL CONCEPTS

The frame of reference used in m theory is  $( \underbrace{r}, \underbrace{\phi} )$ , which is a self consistent development of the plane polar coordinates  $(\underline{r}, \underbrace{\phi})$ . Here:

$$\Gamma_1 = \frac{\Gamma}{m(r)^{1/2}} - (1)$$

where the m  $(r_1)$  function is defined by the infinitesimal line element of the most general spherically symmetric spacetime:

$$ds^{2} = c^{2}dz^{2} = m(r)c^{2}dt^{2} - \frac{dr}{m(r)} - \frac{r}{m(r)} h \phi^{2} - \frac{(1)}{m(r)}$$

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Here  $\tau$  is the proper time and c the speed of light in vacuo, taken to be a universal constant. The hamiltonian of m theory is:

$$H = \chi_m(r_i)mc - mmb - (3)$$

where the potential energy of interaction of a mass m orbiting a mass M is:

$$U = -\frac{m}{r_{1}} - \frac{(4)}{r_{1}} - \frac{r_{1} \cdot r_{1}}{c^{2}} - \frac{(5)}{c^{2}}$$

Here

is the Lorentz factor in m space {1-41}. The lagrangian of m space is:  $J = -mc^{2} \left( m(r_{1}) - \frac{1}{2} \frac{r_{1}}{r_{1}} - \frac{r_{1}}{r_{1}} \right)^{1/2} + nMb - (b)$ 

and fr

It is shown in Note 418(2) and 418(5) that the vector Euler Lagrange equation:

$$\frac{d}{dt}\frac{JJ}{J\dot{t}_{1}} = \frac{JJ}{J\dot{t}_{1}} = \frac{J$$

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is rigorously equivalent to the scalar Euler Lagrange equations:

$$\frac{d}{dt}\frac{\partial I}{\partial i_1} = \frac{\partial I}{\partial r_1} - (8)$$

and

$$\frac{d}{dt} \frac{JL}{J\dot{\phi}} = \frac{JL}{J\phi} - (9)$$

$$\frac{f_{1}}{f_{1}} \cdot \frac{f_{1}}{f_{1}} = \frac{f_{1}}{f_{1}} + f_{1} \frac{\dot{\phi}}{\dot{\phi}} - (10)$$

where:

Eq. ( 9

) gives the conserved angular momentum of m theory:  

$$L = \frac{\partial J}{\partial \dot{\phi}} = \sqrt{mr} \dot{\phi} - (n)$$

$$\frac{dL}{dt} = 0. - (n)$$

so that:

The hamiltonian of m theory is conserved, so:

$$\frac{dH}{dt} = 0 - (13)$$

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The equations of motion of m theory are therefore:

$$\frac{dH}{dt} = 0 - (11+)$$

and

= 0 - (15)

These are rigorously consistent with the equations of motion from the lagrangian of m theory. In the preceding paper (UFT417) it was shown that the hamiltonian, lagrangian and equations of motion give forward and retrograde precession of orbits, shrinking orbits, the possibility of expanding orbits and counter gravitation, superluminal signalling and infinite potential energy from m (r). These are major advances in classical dynamics.

The relativistic kinetic energy of m theory is evaluated from the work integral:

$$W_{12} = T_2 - T_1 = \int_1^d F_1 \cdot dr - (16)$$

from state 1 to state 2. The force in m theory is:

V

$$F_{1} = \frac{dP_{1}}{dt} = \frac{d}{dt} \left( Y_{m} \underline{r}_{1} \right)^{-(17)}$$

 $\chi$  is the Lorentz factor of m theory, first derived in immediately preceding papers. If the initial state is for a particle at rest, and the final state for a particle with velocity  $V_1$ , the relativistic kinetic energy of m theory is: I,  $[v_i \in ]$ ,  $c_i$ ,  $dt = m[v_i d[\delta v_i]$ 

$$W = T = \int \frac{d}{dt} \left( V_{m} \underline{f}_{1} \right) \cdot \underline{f}_{1} dt = m_{0} v_{1} v_{1} (v_{0})$$
by parts:
$$\int_{0}^{N_{1}} v_{1} d(Yv_{1}) = Vv_{1}^{2} - \int_{0}^{N_{1}} Vv_{1} dv_{1} - (19)$$

$$V = \left( m(f_{1}) - \frac{v_{1}}{2} \right)^{-1/2} - (20)$$

Integrating by parts:

where

It follows that:

$$M \int_{0}^{V_{1}} \sqrt{V_{1}} dV_{1} = -mc \left( m(r_{1}) - \frac{V_{1}}{c^{2}} \right)^{1/2} \int_{0}^{V_{1}} - (21)$$

because:

$$\frac{d}{dN_1}\left(m(r_1)-\frac{N_1}{c^2}\right)^{1/2} = -\frac{N_1}{c^2}\left(m(r_1)-\frac{N_1}{c^2}\right)^{1/2}$$

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It has been assumed that:

$$\frac{dm(r_1)}{dv_1} = 0 - (23)$$

The reason for this is that the general spherically symmetric spacetime is:

$$ds^{2} = g_{aa}(a,b)da^{2} + g_{ab}(a,b)(dadb + dbda) + g_{bb}(a,b)db^{2} + r^{2}(a,b)dS^{2} + g_{bb}(a,b)db^{2} + r^{2}(a,b)dS^{2}$$

in spherical polar coordinates. Here a and b are coordinates and (a, b) is an undetermined function. The coordinate function  $\mathcal{K}(\mathcal{K})$  therefore does not depend on the velocity  $\forall_{1}$  of a particle, and Eq. ( 22) follows.  $, 1/2, \sqrt{1}$ 

From Eq. (21):  

$$m \int_{0}^{N_{1}} V_{1} d(X_{1}) = m V_{1}^{2} + m c \left(m(r_{1}) - \frac{V_{1}}{c}\right) |_{0}$$

$$-(25)$$

Therefore the relativistic kinetic energy of m space is:

e relativistic kinetic energy of m space is:  

$$T = \chi m v_{1}^{2} + \frac{mc^{2}}{\chi} - m(r_{1})^{2}mc^{2} - (26)$$

From the generalized Lorentz factor ( 20):

$$v_{i}^{2} = c^{2}\left(w(r_{i}) - \frac{r_{i}}{r_{i}}\right) - (27)$$

$$T = Vmc^{2} \left( m(r_{1})^{2} - \frac{1}{\gamma^{2}} \right) + \frac{mc}{\gamma} - m(r)^{1/2}mc^{2}$$
  
=  $m(r_{1})^{1/2}mc^{2} - m(r_{1})^{1/2}mc^{2} - (28)$ 

From the infinitesimal line element (  $\lambda$  ) the total relativistic energy in m space is:

$$E = m(r_i) \forall m_c^2 - (2q)$$

Therefore the relativistic kinetic energy in m space is:  $T = E - n(r) mc^{2} - (30)$ 

The Einstein energy equation in m space is found from the relativistic linear

momentum in m space:  

$$P_{1} = \sqrt{m} \sqrt{1} = \sqrt{m} \frac{r_{1}}{1} - (31)$$

From Eq. (31):  

$$P_1 c = \gamma m \gamma_1 c = \gamma m c + \frac{\gamma_1}{c^3} - (32)$$
  
where  
 $\gamma_1 = m(r_1) - \frac{1}{c^3} - (33)$ 

from the Lorentz factor of m space, Eq. (
$$20$$
). It follows that:  
 $P_{1}c^{2} = 3m^{2}c^{4}\left(m(1)^{2} - \frac{1}{\sqrt{2}}\right) - (34)$ 

and:

$$E^{2} = P_{1}^{2} C^{+} E^{2}_{0} - (35)$$

which is the Einstein energy equation in m space, Q. E. D. Here

$$E_0 = nc$$
  $-(36)$ 

is the rest energy in m space. Defining the four momentum:

$$P_{1}^{\mathcal{M}} = \left(\frac{E}{c}, \frac{P_{1}}{c}\right)^{-(37)}$$

the Einstein energy equation in m space becomes:

$$P_{1}^{\mu}P_{1\mu} = mc. - (38)$$

 $\langle \rangle \rangle$ 

This is an equation of ECE and ECE2 generally covariant unified field theory and can be quantized to a wave equation and related to the tetrad postulate of Cartan geometry.

The classical dynamics of m space can be developed using the force equation:

$$\overline{F}_{1} = \frac{d \overline{f}_{1}}{d t} = -\left(\frac{nc}{2} \chi \frac{d n(r_{1})}{d r_{1}} - \frac{n M f}{r_{1}}\right) \frac{e}{-(3q)}$$

which is the direct result of the Euler Lagrange equation (7) with the m space lagrangian (**b**). This contains the new force term:

$$F = -\frac{mc}{2} \frac{\chi dm(r_i) g}{dr_i} - (40)$$

which in UFT417 was defined as the vacuum force. In the  $(r, \phi)$  coordinate system:  $F = mc \left( \frac{Vm(r)}{rdm(r)} - 2m(r) \right) \frac{dm(r)}{dr} + r - (41)$ 

and the vacuum force goes to infinity when:

$$r dn(r) = dn(r) - (4d)$$

Using the work theorem:  

$$T_2 - T_1 = U_1 - U_2 = -mc \frac{1}{2} \int \frac{dm(r_1)}{dr_1} dr_1 - (43)$$

as usual in classical dynamics. Here:

ics. Here:  

$$\Delta U = \overline{U_1} - \overline{U_2} - (44)$$

is the change in potential energy. In flat spacetime:

In flat spacetime:  

$$\frac{dn(r_{1})}{dr_{1}} = 0 - (45)$$

$$\int \sqrt{r} = 0 - (47)$$

so:

From Eq. (
$$L_3$$
), total energy is conserved:  
 $T_1 + U_1 = T_2 + U_2$ .  $-(L_7)$ 

The change in potential energy due to the vacuum force is therefore:  

$$MI = -\frac{mc}{2} \int_{1}^{2} \chi dm(r_{1}) = -\frac{mc}{2} \int_{1}^{2} -(48)$$

If it is assumed that the initial state is that of a particle at rest:

$$V_1 = 0 - (49)$$

then:

$$\Delta U = -mc^{2} \left( \left( m(r_{1})^{2} - \frac{v_{1}}{c} \right)^{1/2} - m(r_{1})^{1/2} - \frac{v_{1}}{c} \right)^{1/2} - (50)^{1/2} \right)$$

Note carefully that this is non-zero if and only if:

$$\frac{dm(i)}{ds_{1}} \neq 0. - (s_{1})$$

$$\frac{dm(i)}{ds_{1}} \neq 0. - (s_{2})$$

In the classical limit:

the change in potential energy due to 
$$m(r_1)$$
 is:  

$$M = -mc^2 m(r_1)^{1/2} \left( \left( \frac{1-v_1}{n(r_1)c^2} \right)^{1/2} - 1 \right) \rightarrow \frac{1}{2} \frac{mv_1}{n(r_1)^{1/2}}$$

$$-(53)$$

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and in the limit of flat spacetime:

$$T = \Delta U \rightarrow \frac{1}{2} mv - (54)$$

in the usual (r,  $\phi$ ) coordinate system.

The velocity v can be thought of as that of a vacuum particle, and is imparted to

the material particle of mass m by the fact that:  $dm(s_1)/ds_1 \neq 0 - (55)$ Integrating by parts as in Note 418(8):  $\int \sqrt{dm(s_1)} ds_1 = \sqrt{m(s_1)} - \int m(s_1) dx_1 - (56)$ The function  $dm(s_1)/ds_1$  becomes infinite at:  $r dm(s_1) = 2m(s_1) - (57)$ 

Self consistently from Eq. (5), in the flat spacetime limit of:

$$\frac{dm(r_i)}{dx_i} = 0 - (58)$$

X -> 00. - (59)

it follows that:

From Eq. (56), the condition (59) can be satisfied by:  

$$n(r_1) = \sqrt{2} / c^2 - (60)$$
i. e.  

$$n(r) = \sqrt{2} \cdot \frac{1}{c^2} - (61)$$

 $m(r) = \frac{x}{c}.$ 

In the ( $(, \phi)$ ) coordinate system this means:

## 3. NUMERICAL ANALYSIS AND GRAPHICS

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